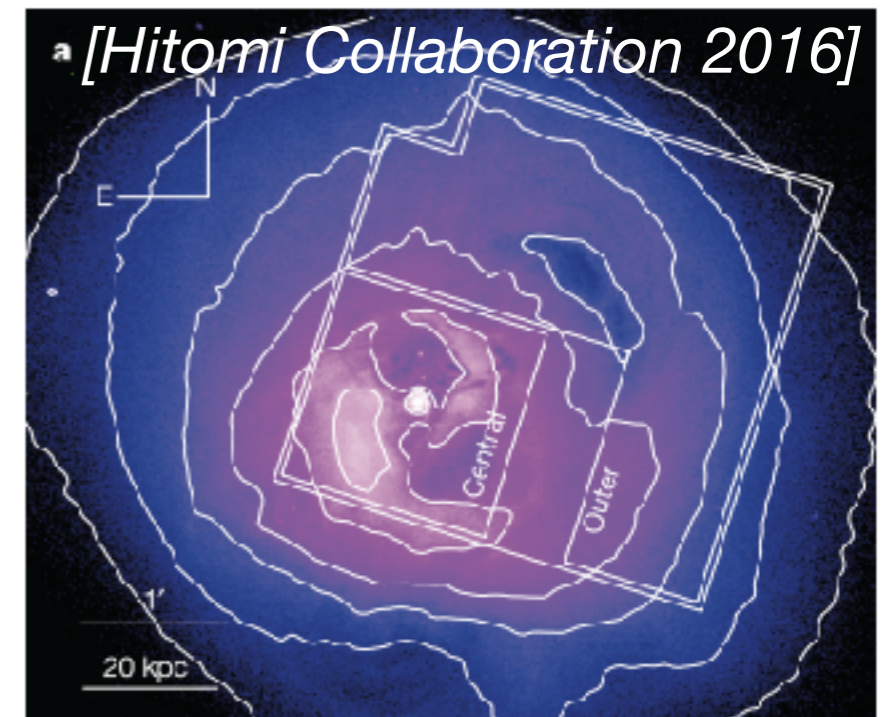
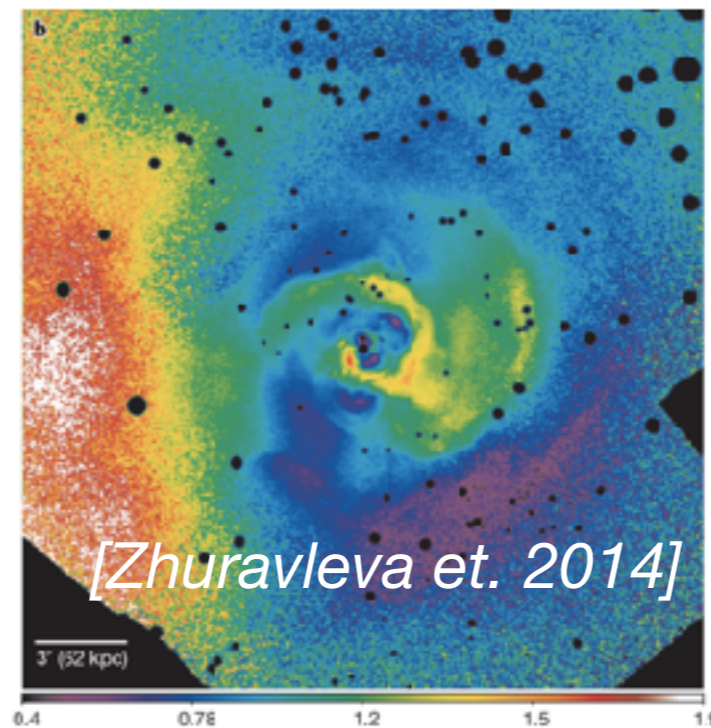
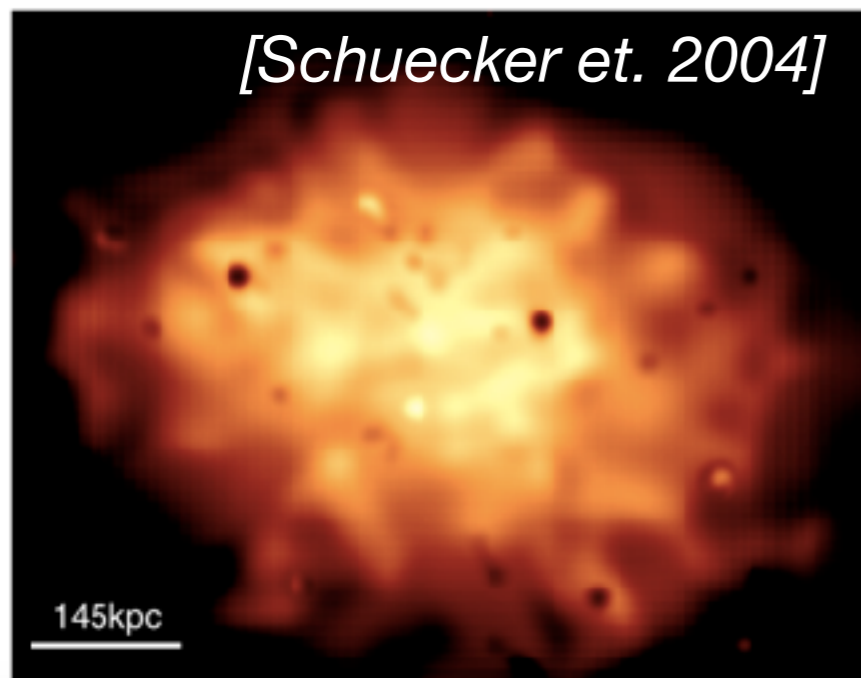


Turbulence in the ICM

Prateek Sharma, IISc Bangalore

ICM workshop, Garching, 09/09/2018



Systematics have really interesting physics!

Why study ICM turbulence?

- Turbulent support biases cluster mass estimates
- Turbulent heating in cool cores
- Constraints on ICM (high- β) plasma transport
- Implications on thermal instability & condensation, extrapolation to puzzling CGM observations (turbulent BLs, cloudlets, $t_{\text{cool}}/t_{\text{eddy}}, \dots$)
- Convert SB (X-ray maps) & projected pressure fluctuations (resolved SZ maps) to v_t

Idealized simulations

Turbulence in the intracluster medium: simulations, observables & thermodynamics

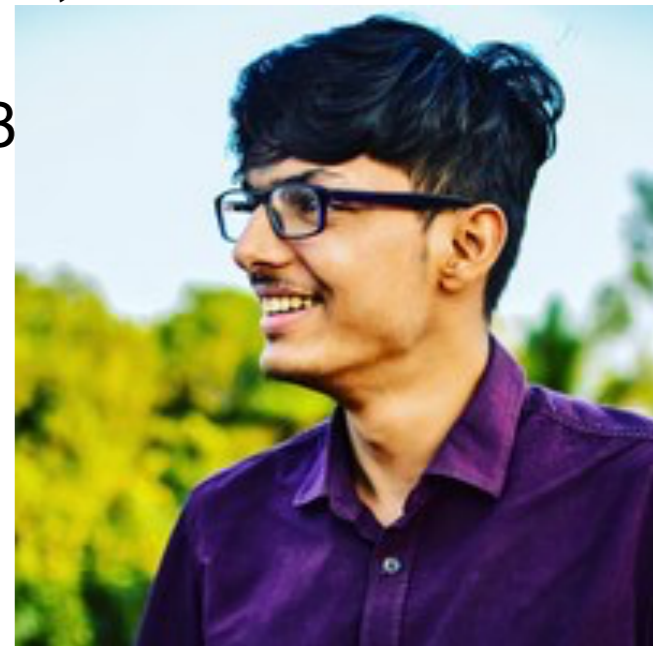
arXiv:1810.00018

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Isotropic, homogeneous turbulence

To test scalings of density/pressure fluctuations with rms Mach number

ρ , P , v Power spectra: ρ/P spectra smaller by a factor of M^2 as compared to v

Turbulence with heating & cooling

TI leads to much larger density fluctuations

If cold gas condenses out of ICM, $t_{\text{cool}} < t_{\text{mix}}$

Constraints v_t & turbulent heating to be small

Idealized simulations

Turbulence in the intracluster medium: simulations, observables & thermodynamics

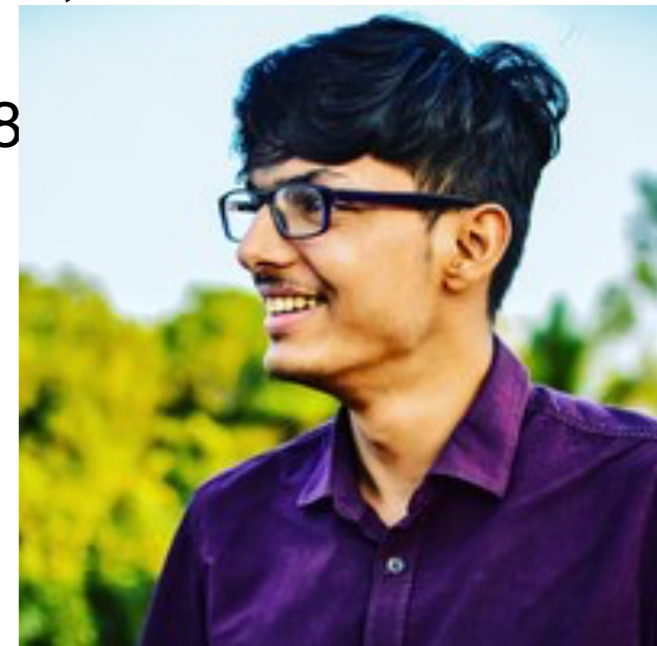
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Isotropic, homogeneous turbulence

Caveats:

ICM is gravitationally stratified
Turbulence can excite gravity waves

$$\delta\rho_k/\rho \sim (v_k/c_s)/\sqrt{3}$$

But Kolmogorov at small scales!

Turbulence with heating & cooling

Caveats:

Cold gas may not condense from hot ICM

E.g., can be uplifted from the BCG

In this case, t_{cool} can be $\gg t_{\text{mix}}$

& turbulent heating can balance cooling

Turbulence simulations

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = 0,$$

$$\frac{\partial(\rho \mathbf{v})}{\partial t} + \nabla \cdot (\rho \mathbf{v} \otimes \mathbf{v}) + \nabla P = \mathbf{F},$$

Solenoidal forcing at selected ks



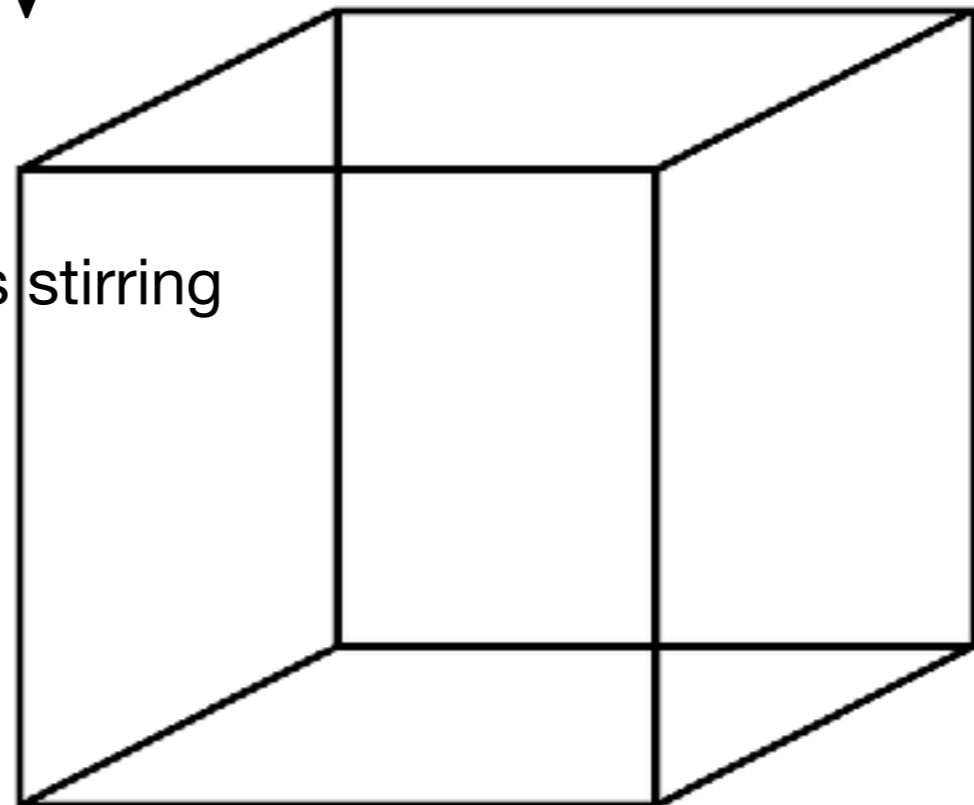
Power input by turbulent forcing heats the gas



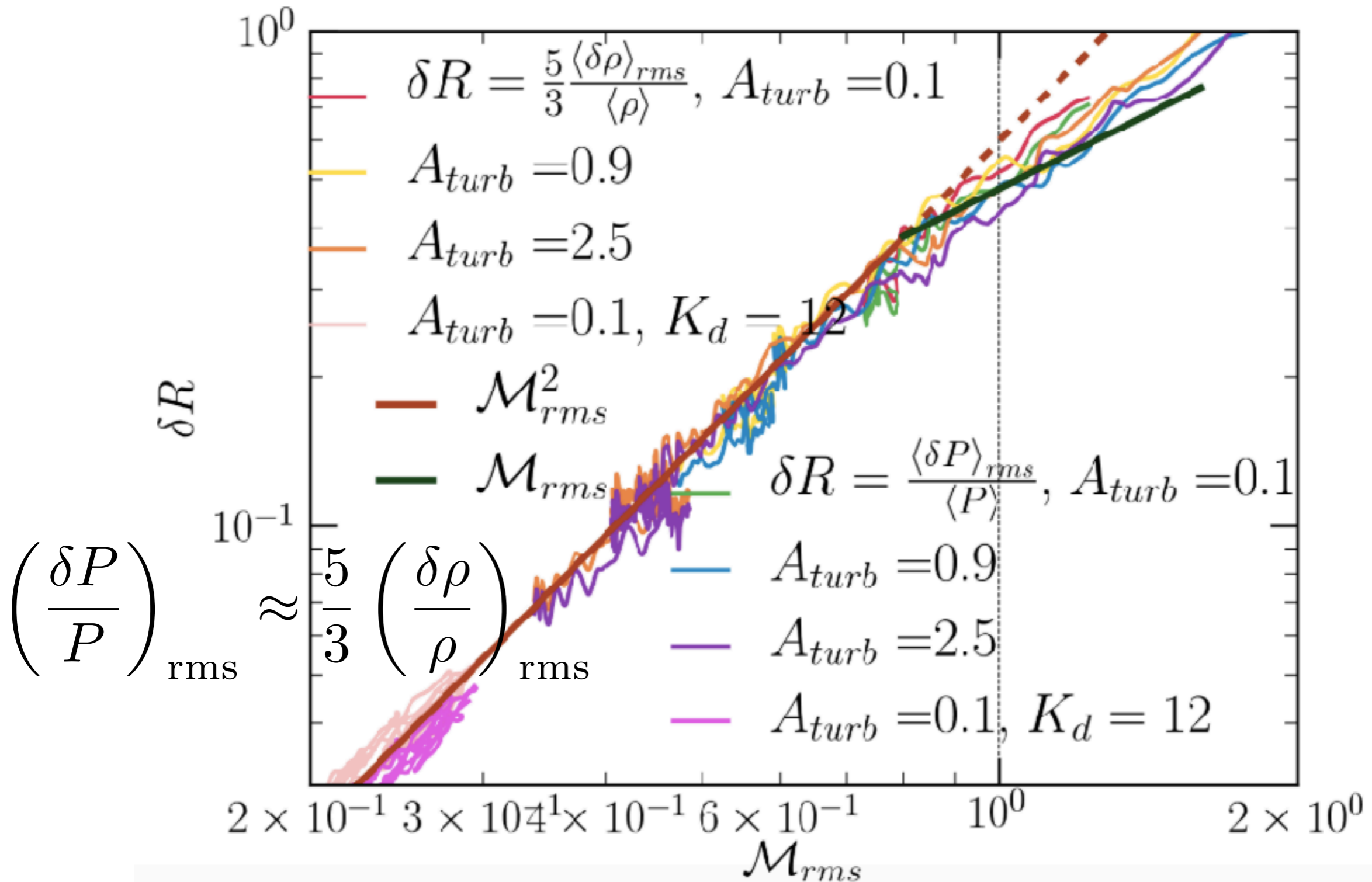
$$\frac{\partial E}{\partial t} + \nabla \cdot ((E + P)\mathbf{v}) = \mathbf{F} \cdot \mathbf{v}$$

3-D periodic box, isotropic/homogeneous stirring

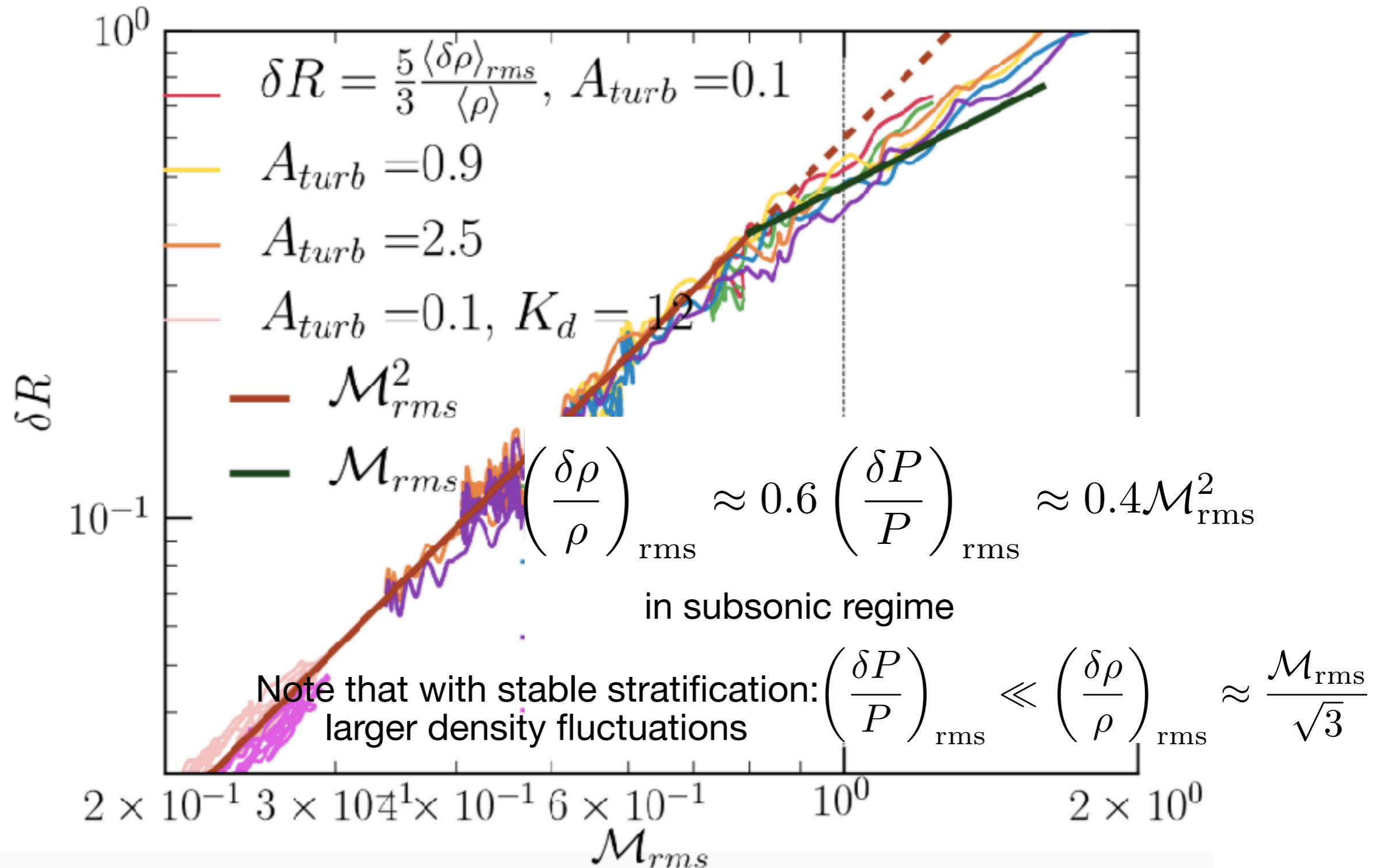
Using PLUTO code
256³, 512³ simulations



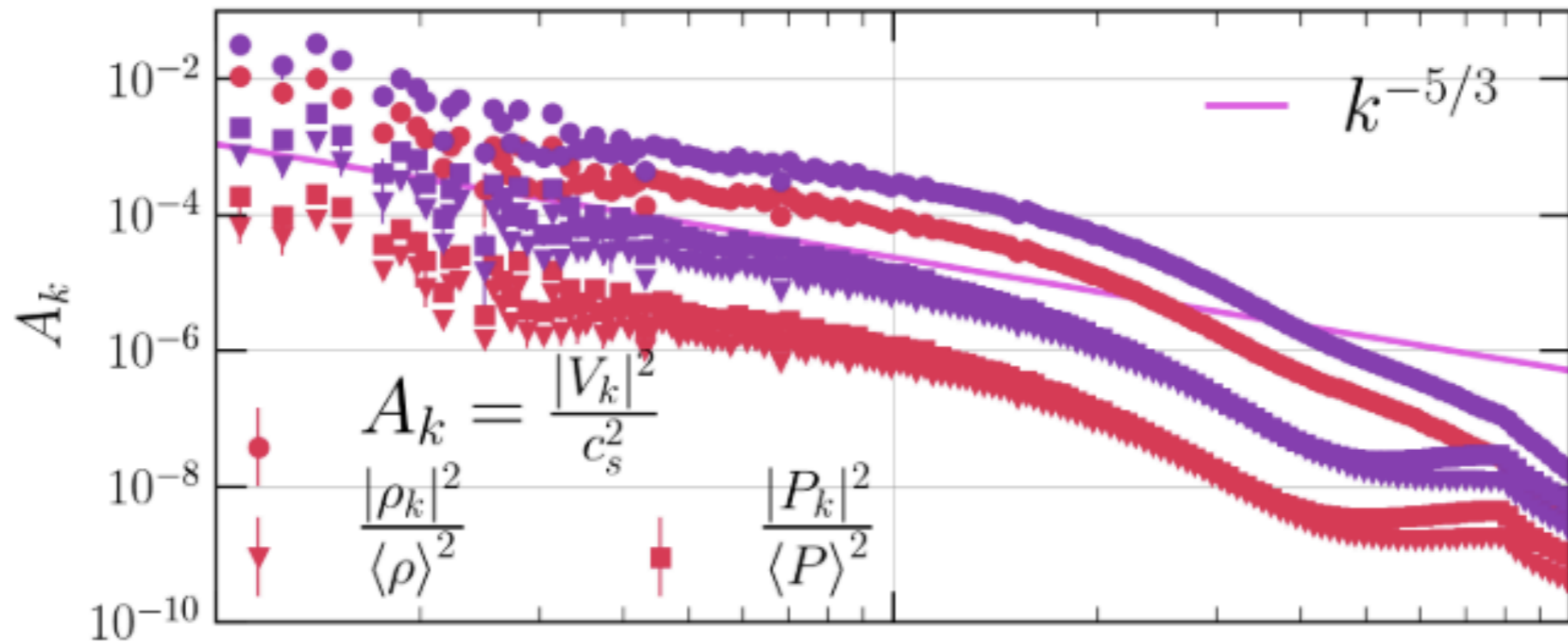
ρ/P fluctuations vs M



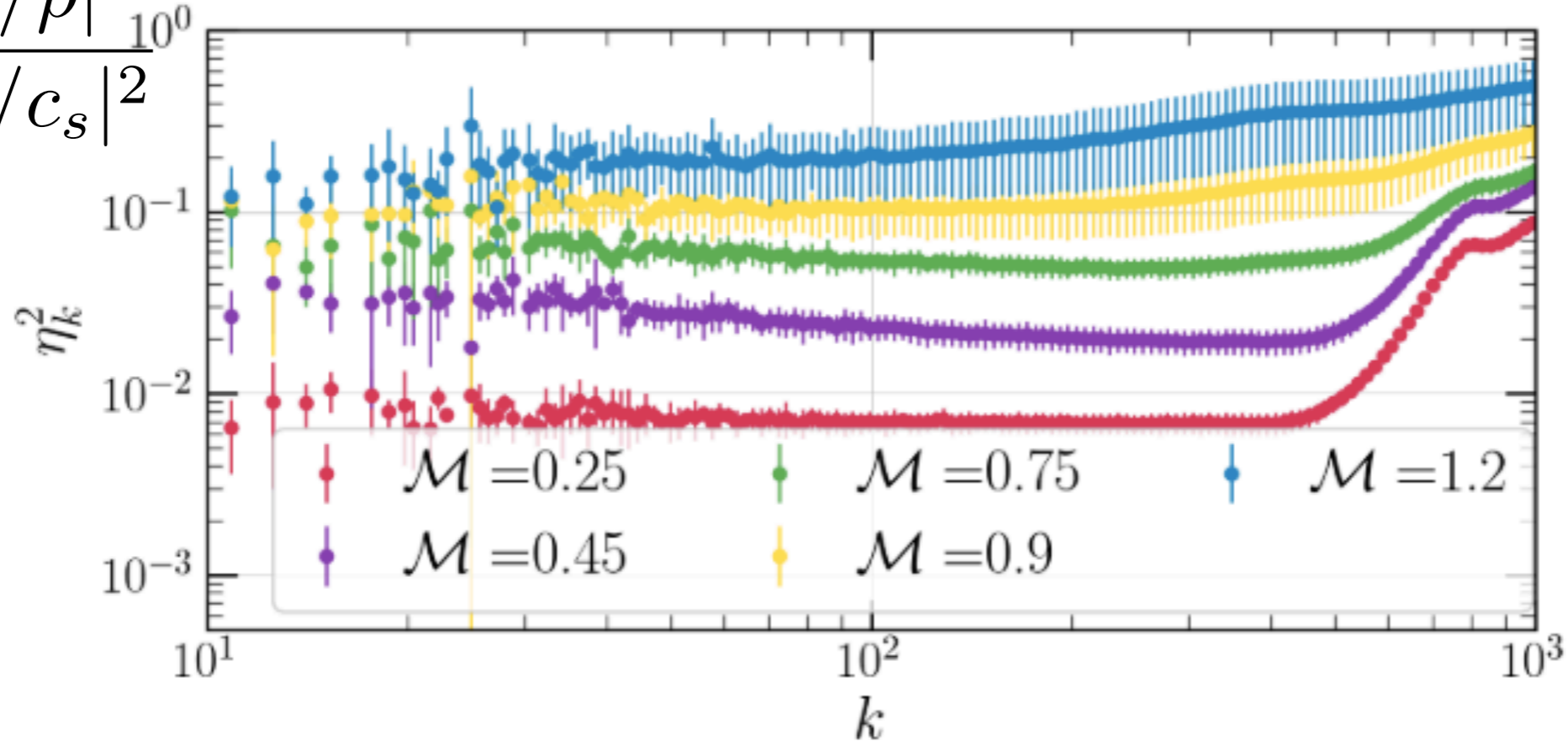
ρ/P fluctuations vs M



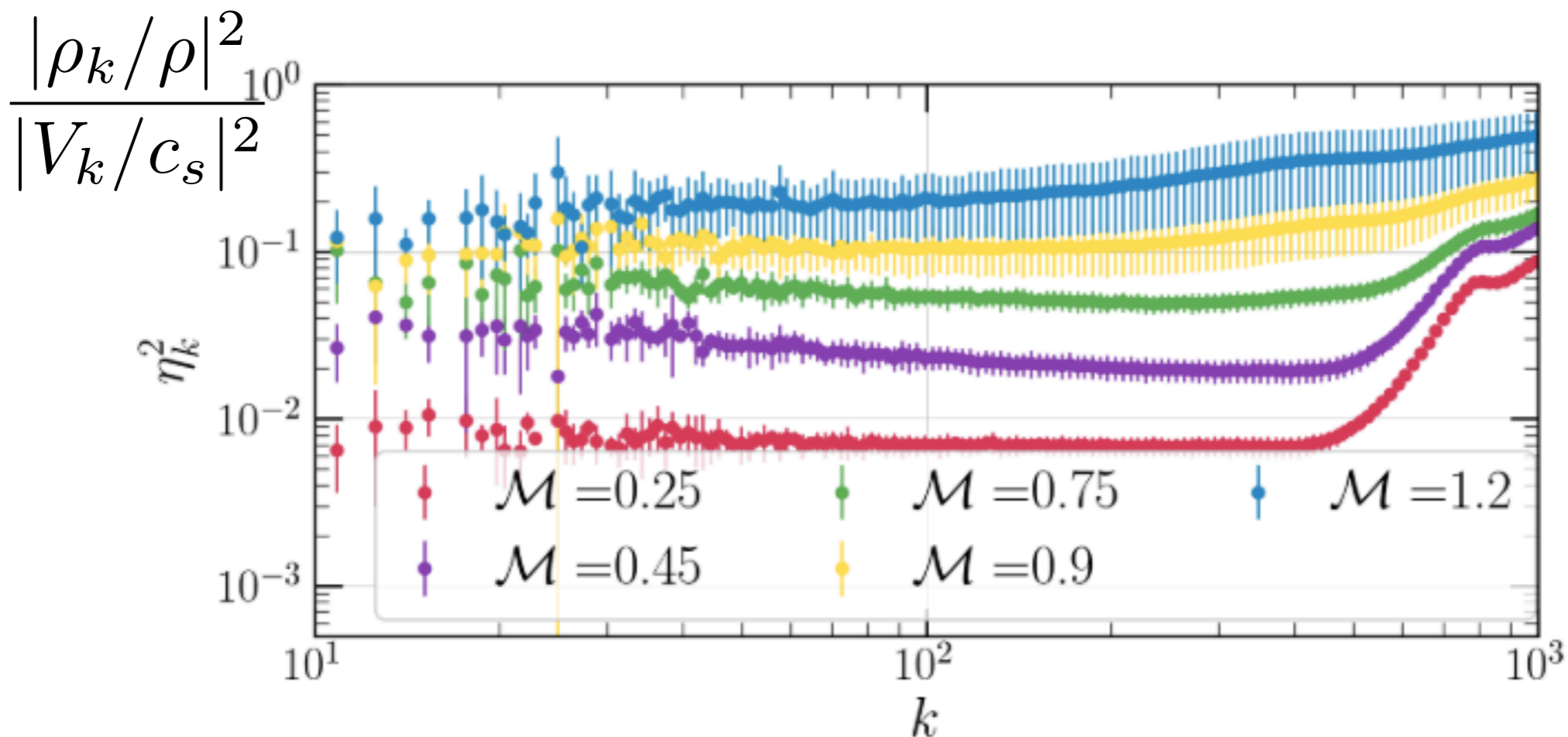
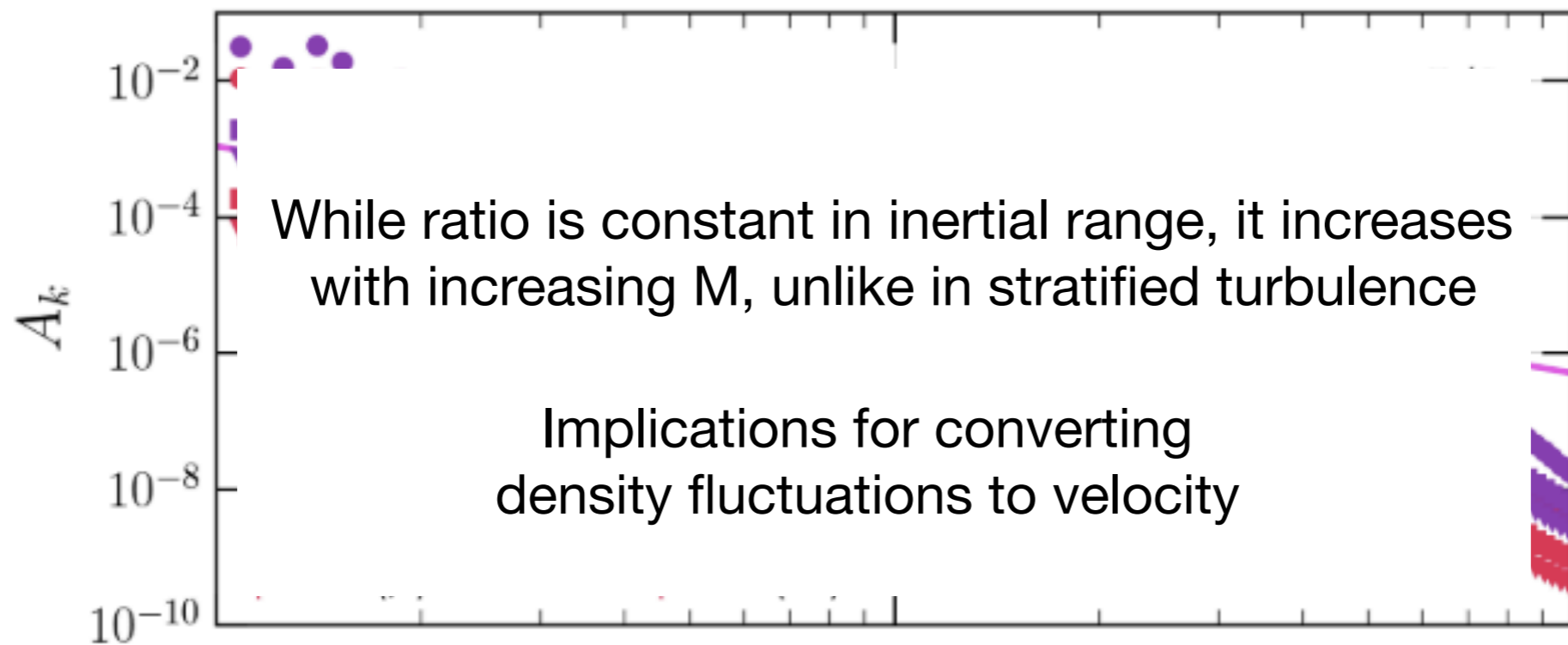
Power spectra



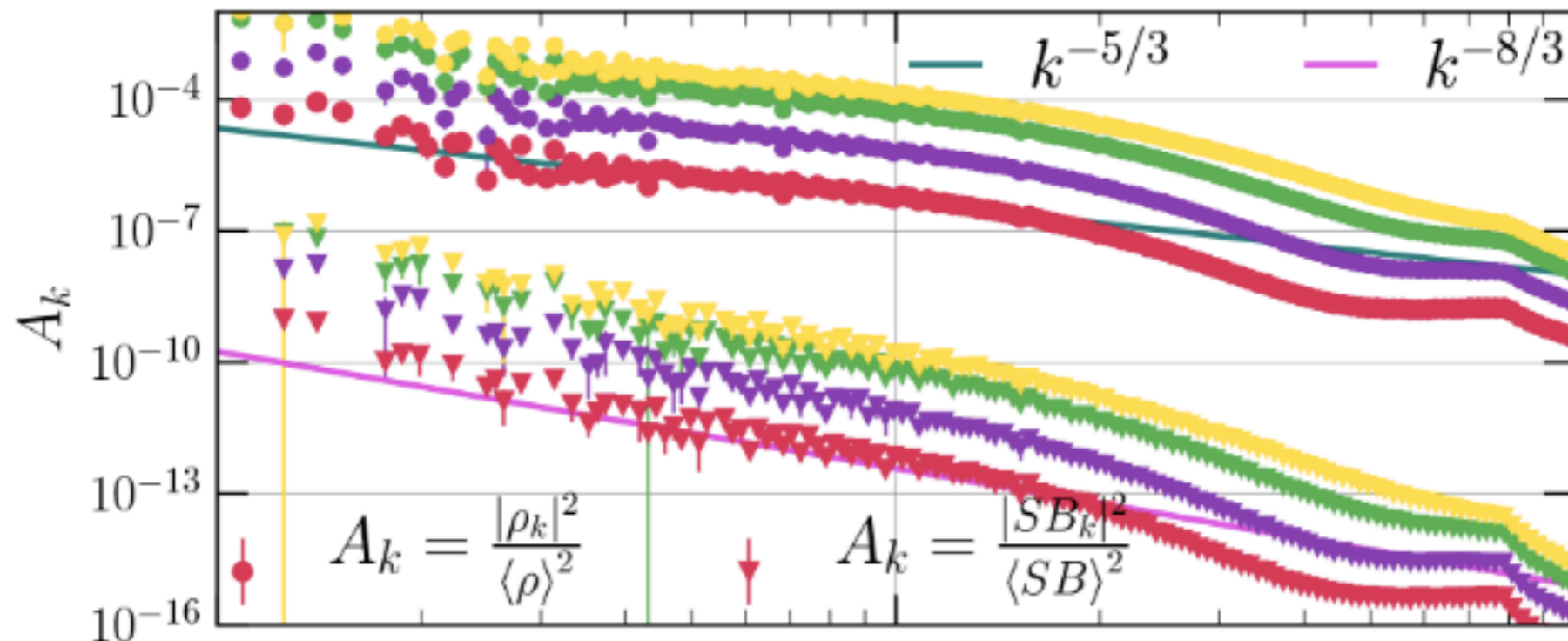
$$\frac{|\rho_k/\rho|^2}{|V_k/c_s|^2}$$



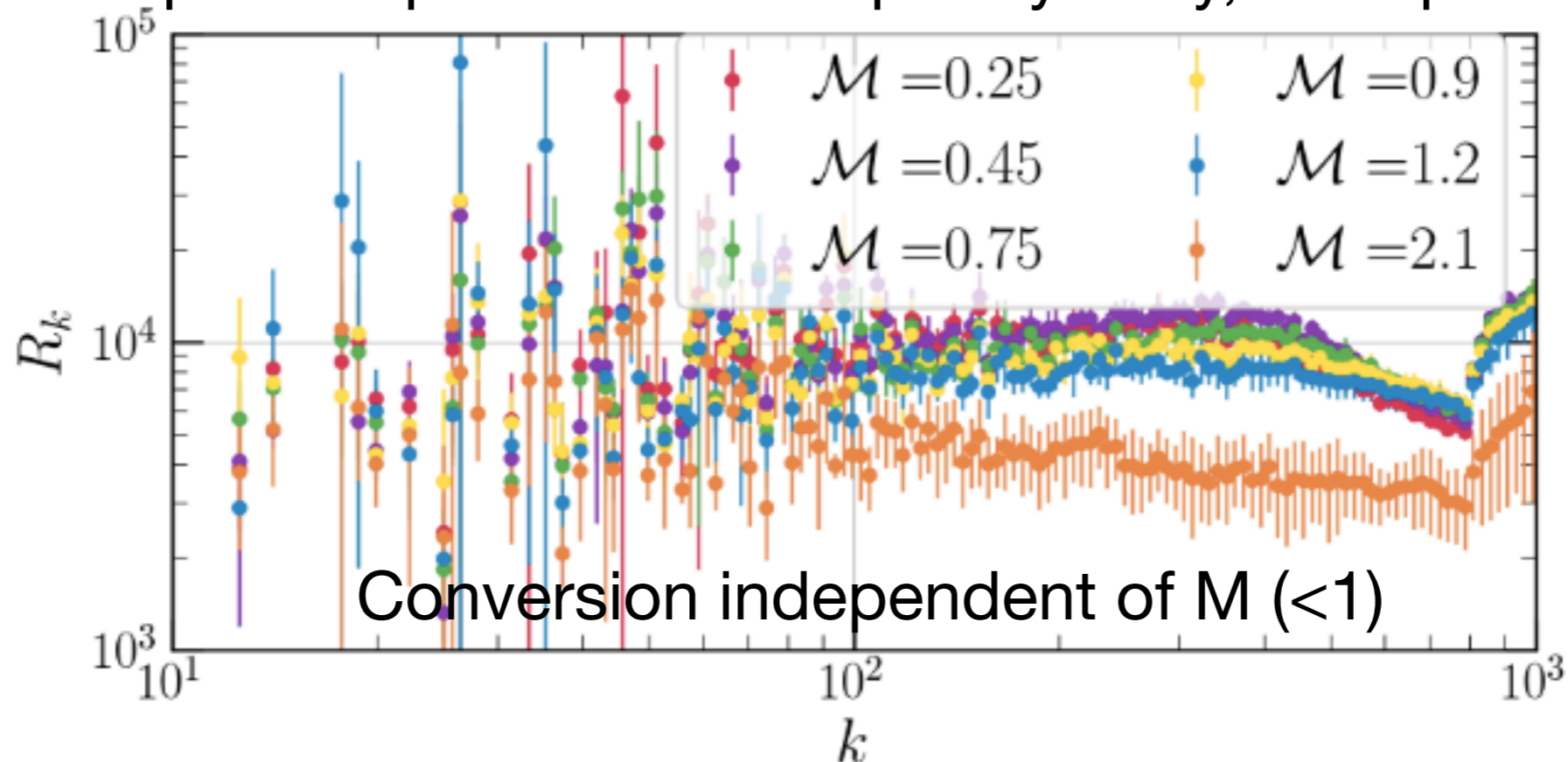
Power spectra



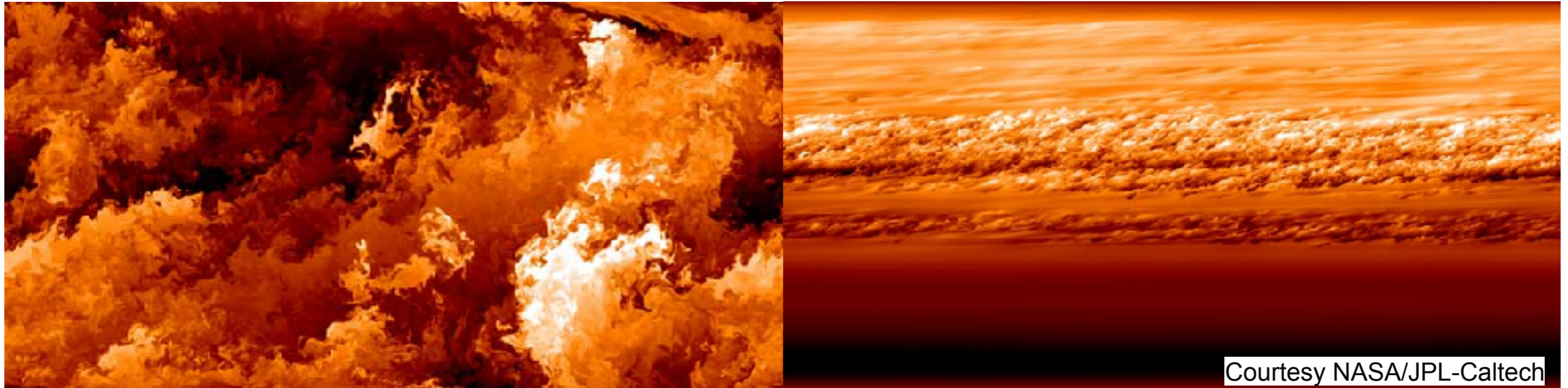
ρ /SB spectra



SB power spectrum is steeper by unity, as expected!



Stably stratified turbulence



$$\text{Ri}(l) = \text{Fr}^{-2}(l) = \frac{N^2}{v^2(l)/l^2} = \frac{\frac{g}{\gamma} \frac{d}{d \ln r} \ln(p/\rho^\gamma)}{v^2(l)/l}$$

Spiral structures in ICM!

Anisotropy crucial for $\text{Ri} \gg 1$, not so much for the ICM that has a shallow entropy profile

$$g \frac{d}{d \ln r} \ln T \text{ for magnetized plasmas, which is even shallower}$$

Turbulence dominates at small scales.

Caveat for the present work: background stratification & gravity are ignored
Background gravity also suppresses multiphase condensation due to local TI

Turbulence, heating & cooling

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = 0,$$

Solenoidal forcing at selected ks

$$\frac{\partial(\rho \mathbf{v})}{\partial t} + \nabla \cdot (\rho \mathbf{v} \otimes \mathbf{v}) + \nabla P = \mathbf{F},$$

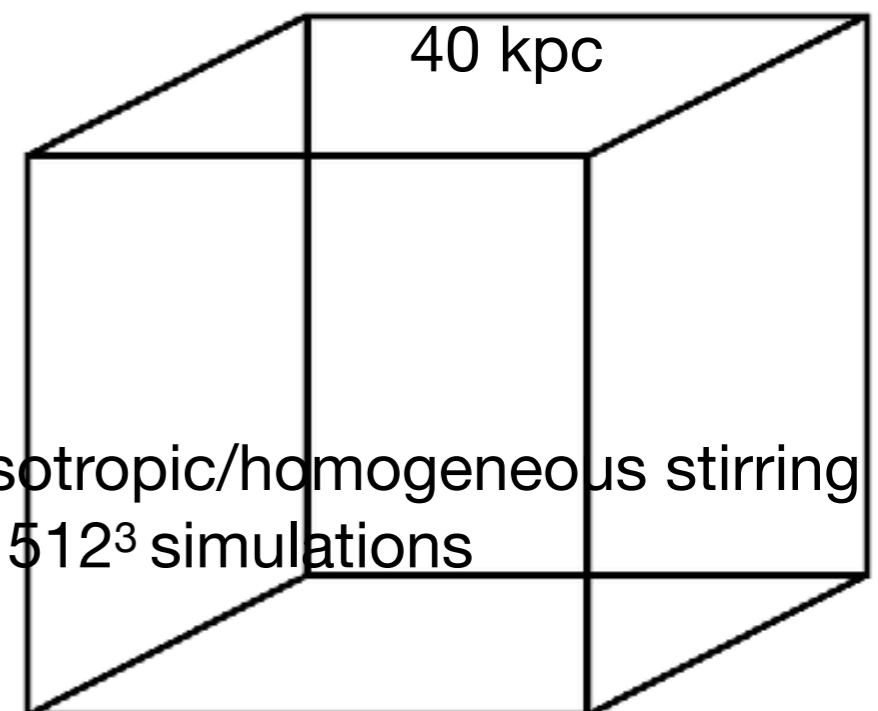
Power input by turbulent forcing = $f_{\text{turb}} \langle \mathcal{L} \rangle$

$$\frac{\partial E}{\partial t} + \nabla \cdot ((E + P)\mathbf{v}) = \mathbf{F} \cdot \mathbf{v} + Q - \mathcal{L} \quad \equiv n_e n_i \Lambda [T] \quad \text{floor } T=10^6 \text{ K}$$

$(1 - f_{\text{turb}}) \langle \mathcal{L} \rangle$
added as thermal energy

Imposed observed thermal balance

3-D periodic box, isotropic/homogeneous stirring
256³, 512³ simulations



K41 Estimates

thermal balance, driving at large scales

$$\dot{E}_{\text{turb}} \sim \rho v_l^2 / t_{\text{mix},l} \approx \rho v_l^3 / l \approx f_{\text{turb}} \dot{E}_{\text{cool}} = f_{\text{turb}} U / t_{\text{cool}}$$

$$t_{\text{cool}} / t_{\text{mix},L} \approx f_{\text{turb}} U / 2K \sim f_{\text{turb}} \mathcal{M}_{\text{rms}}^{-2}$$

If cold gas is to condense out of the hot phase, $t_{\text{cool}} \lesssim t_{\text{mix}}$

or, $\mathcal{M} \gtrsim 1$ for $f_{\text{turb}} \sim 1$

For driving at small scales, an equivalent estimate is

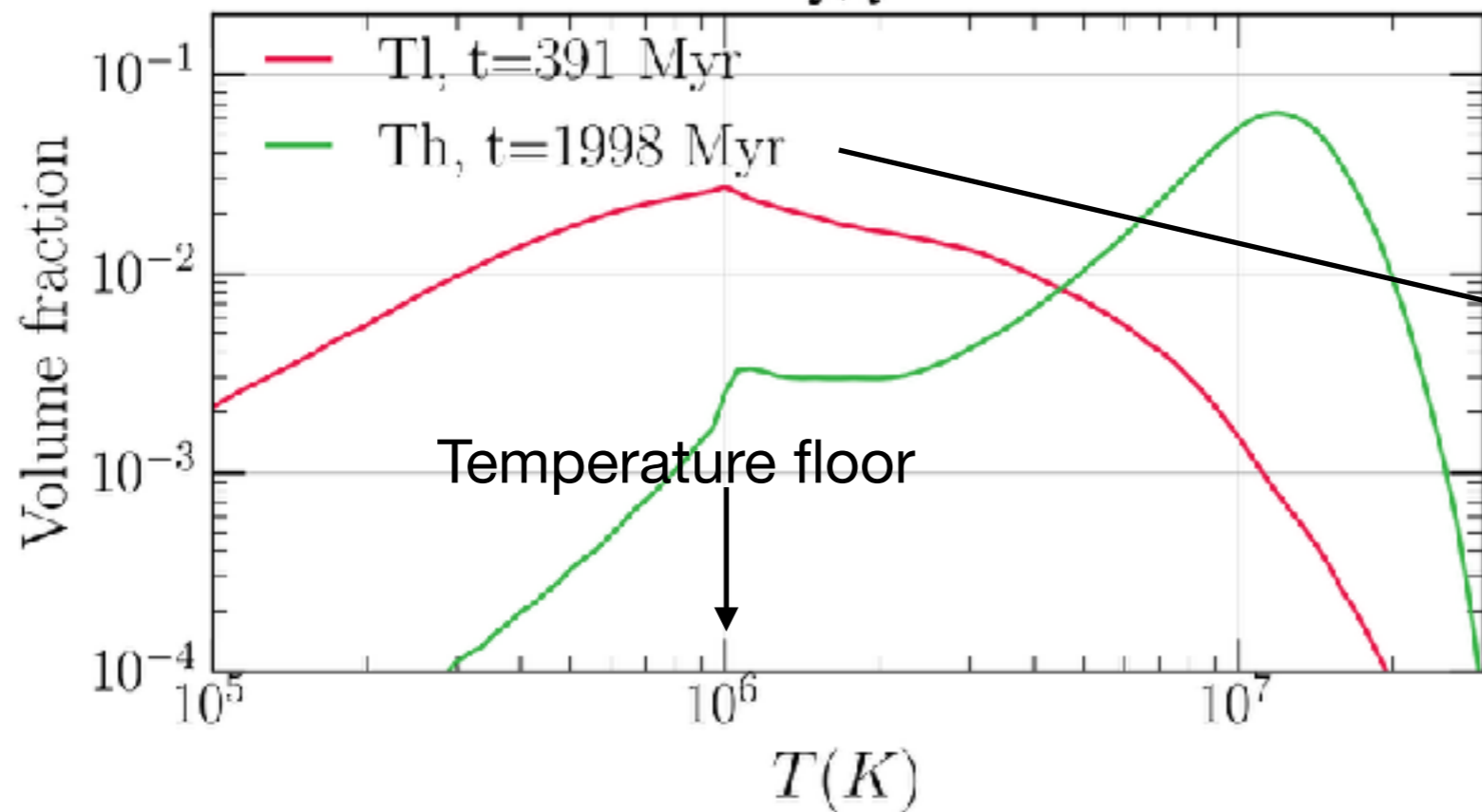
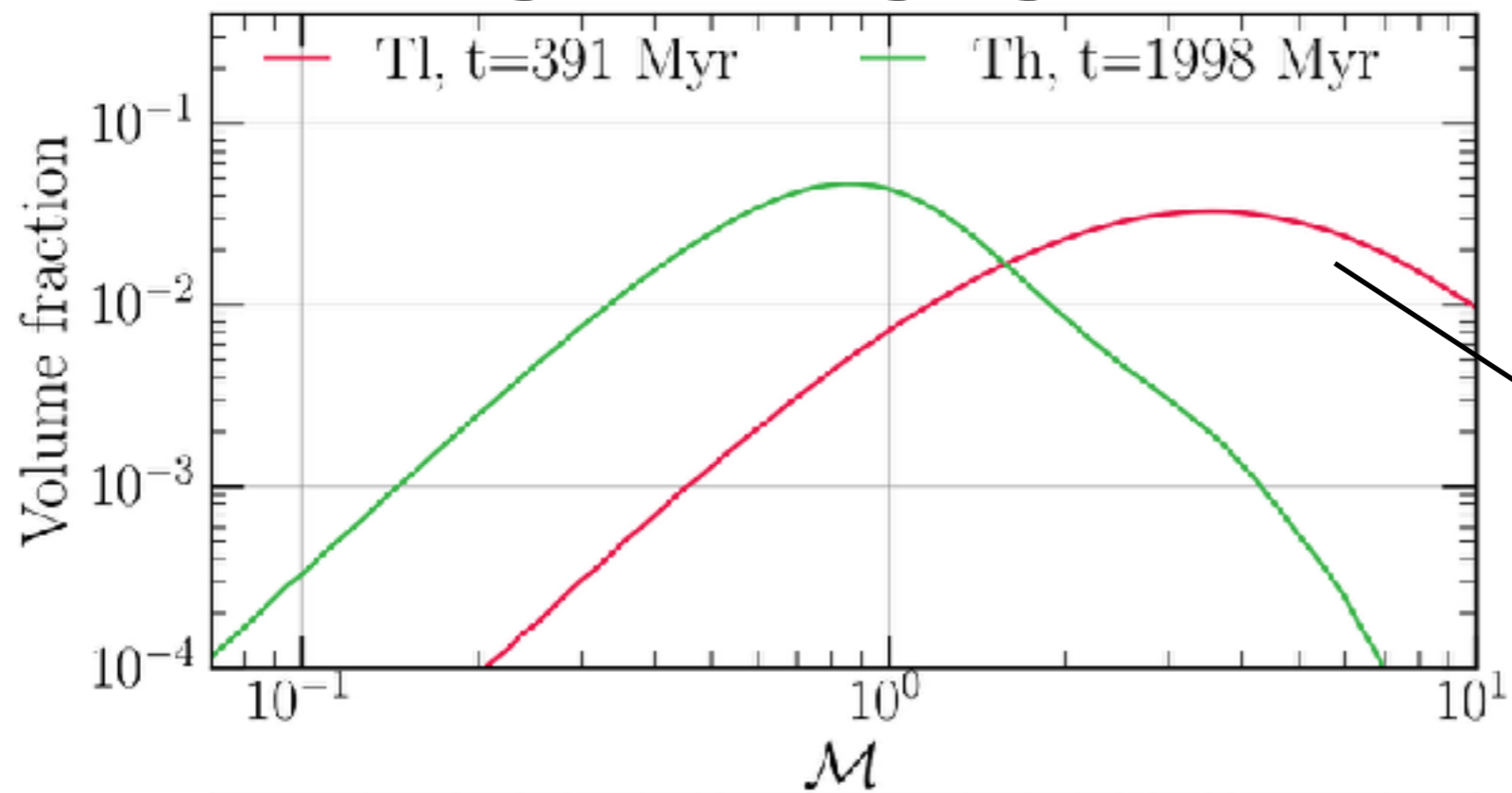
$$f_{\text{turb}} (L/l)^2 \mathcal{M}_{\text{rms}}^{-2} < 1 \text{ for } l > L$$

A smaller/reasonable \mathcal{M}_{rms} for small scale driving & smaller f_{turb}

$t_{\text{cool}} / t_{\text{mix}} \gg 1$ Ok if most cold gas does not condense out but say is uplifted

$k = \sqrt{2}(2\pi/L)$ $k = 12(2\pi/L)$

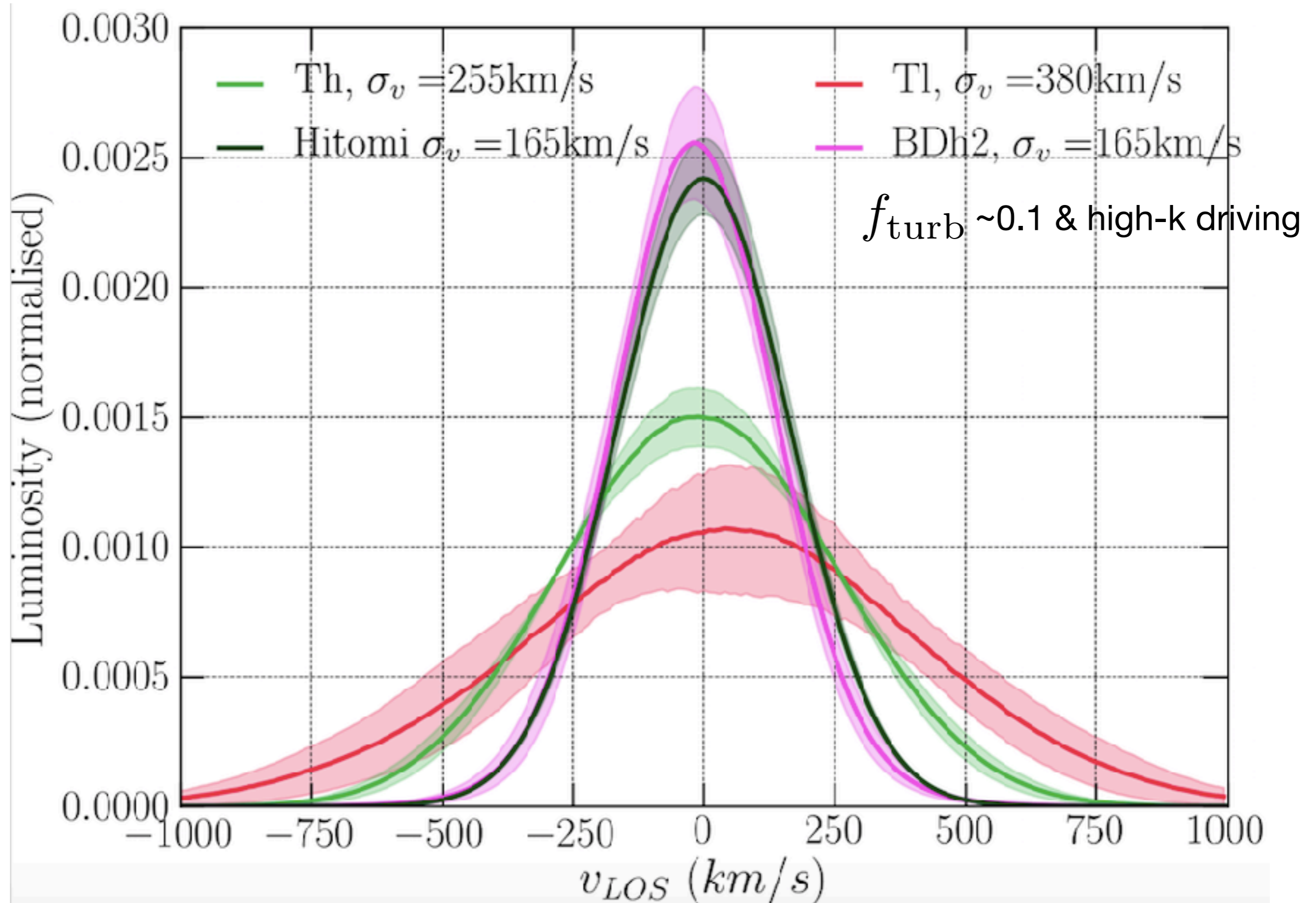
LS & SS driving



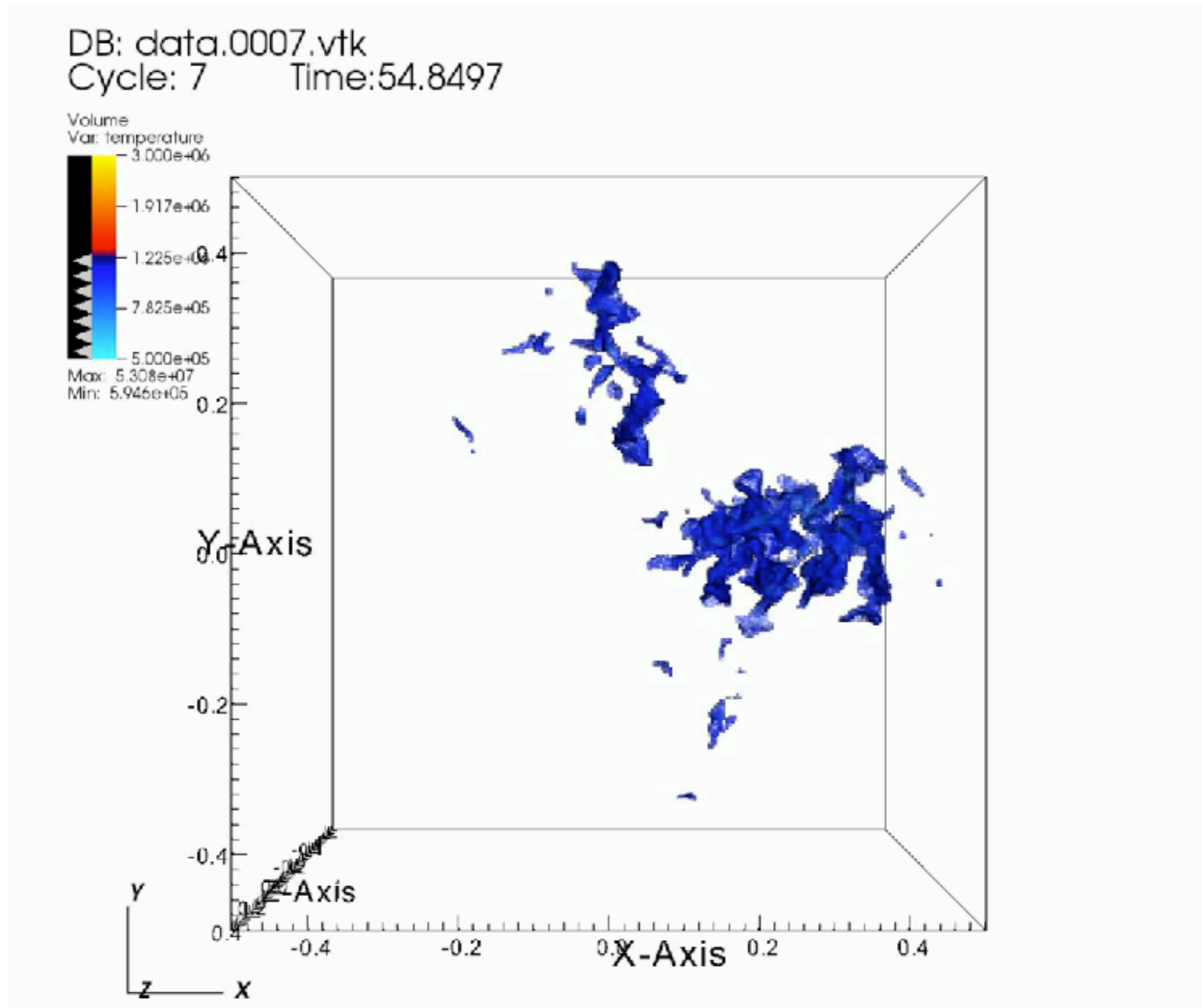
Condensation takes too long for SS driving, fast mixing

Very large velocities for both!

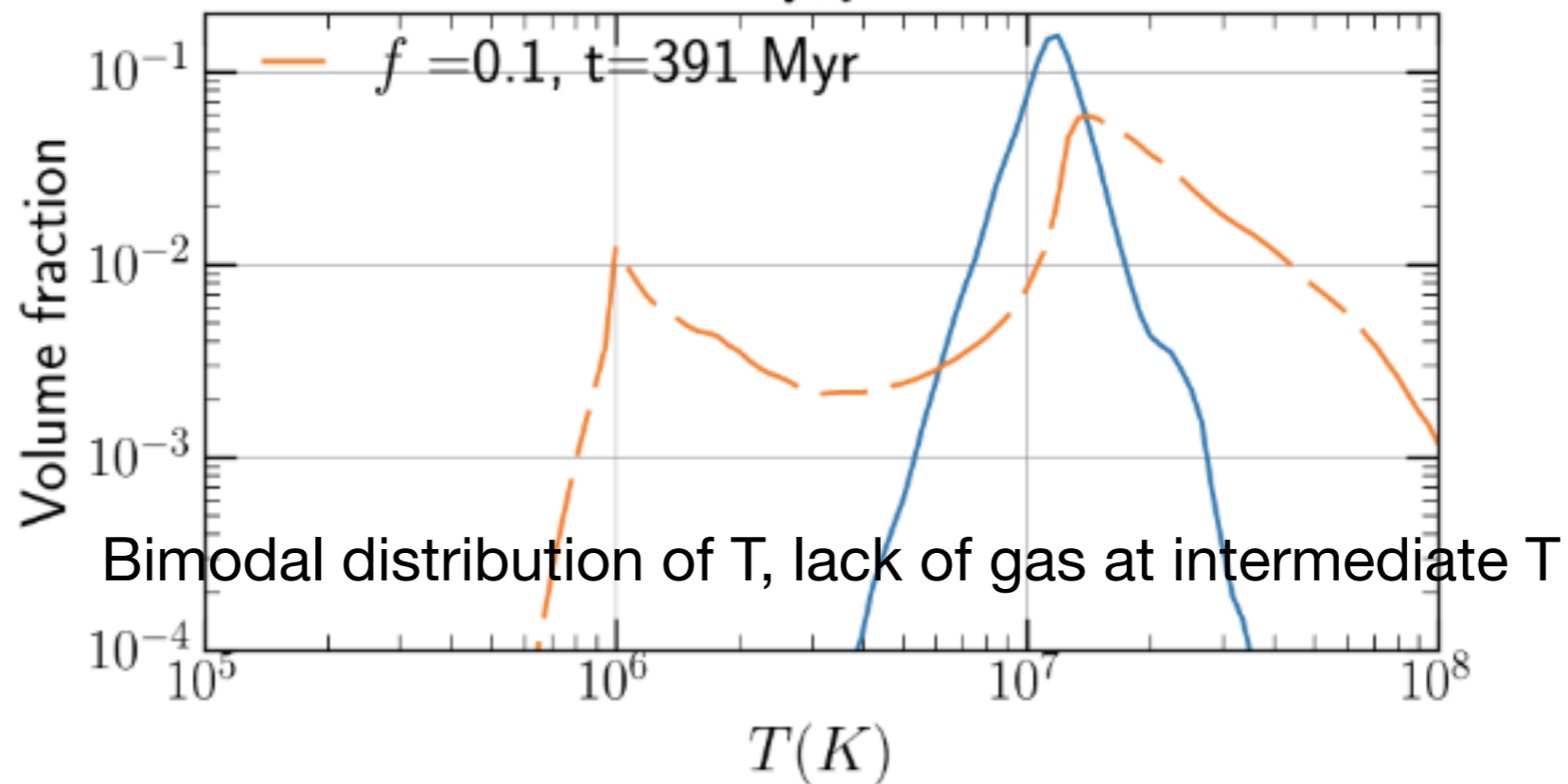
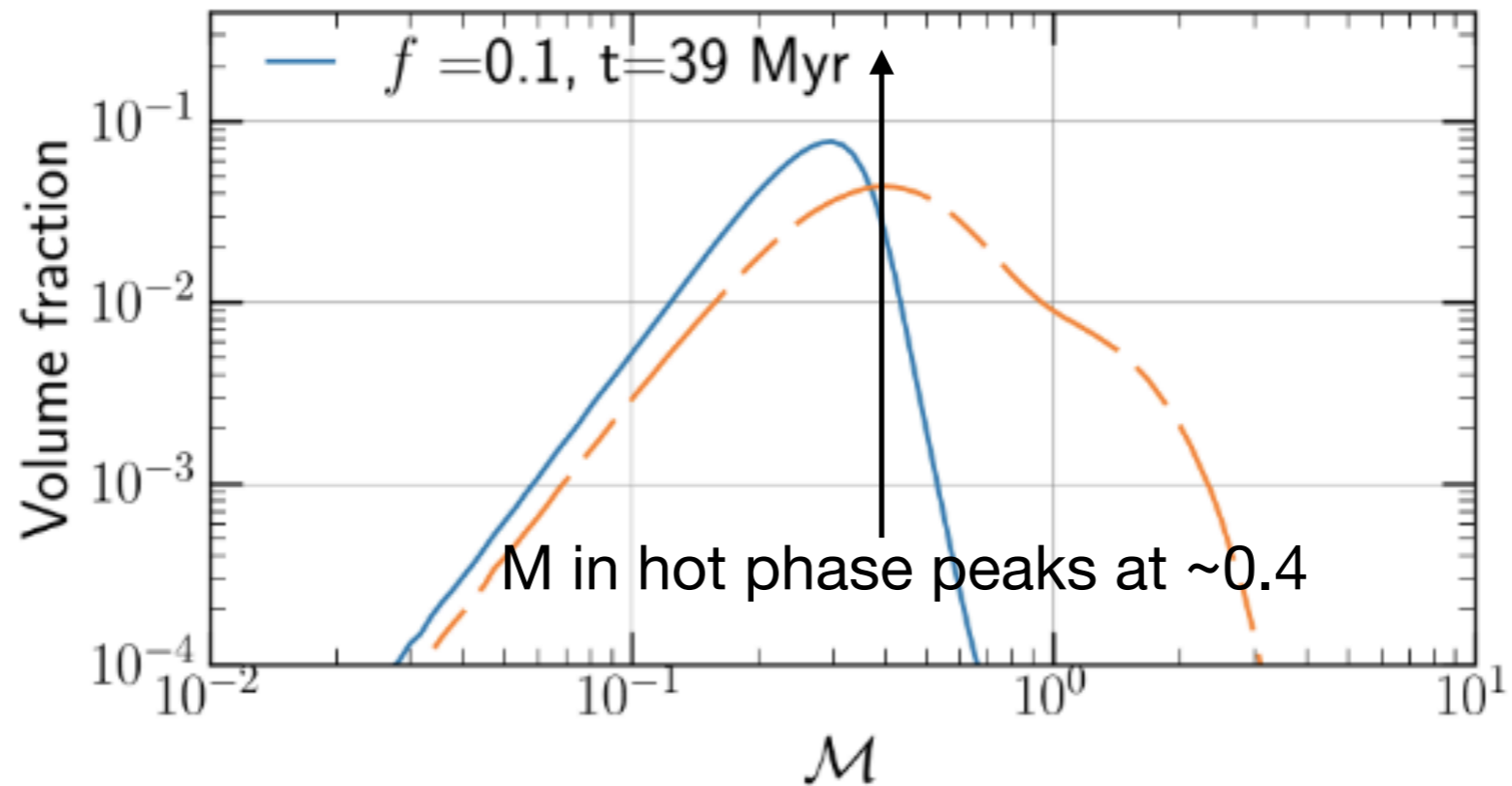
Adjusting f_{turb} to match Hitomi



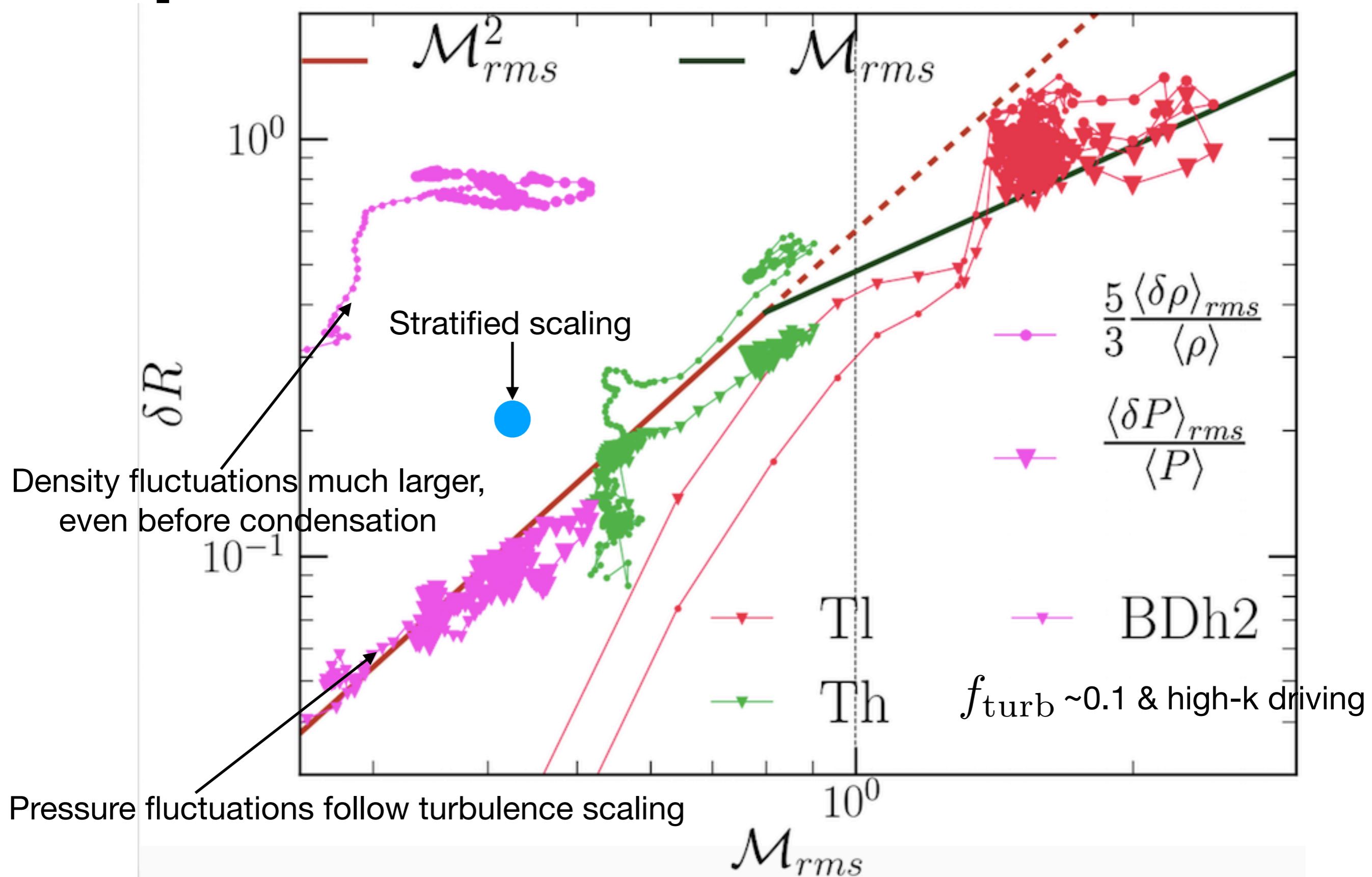
Condensation for $f_{\text{turb}} \sim 0.1$



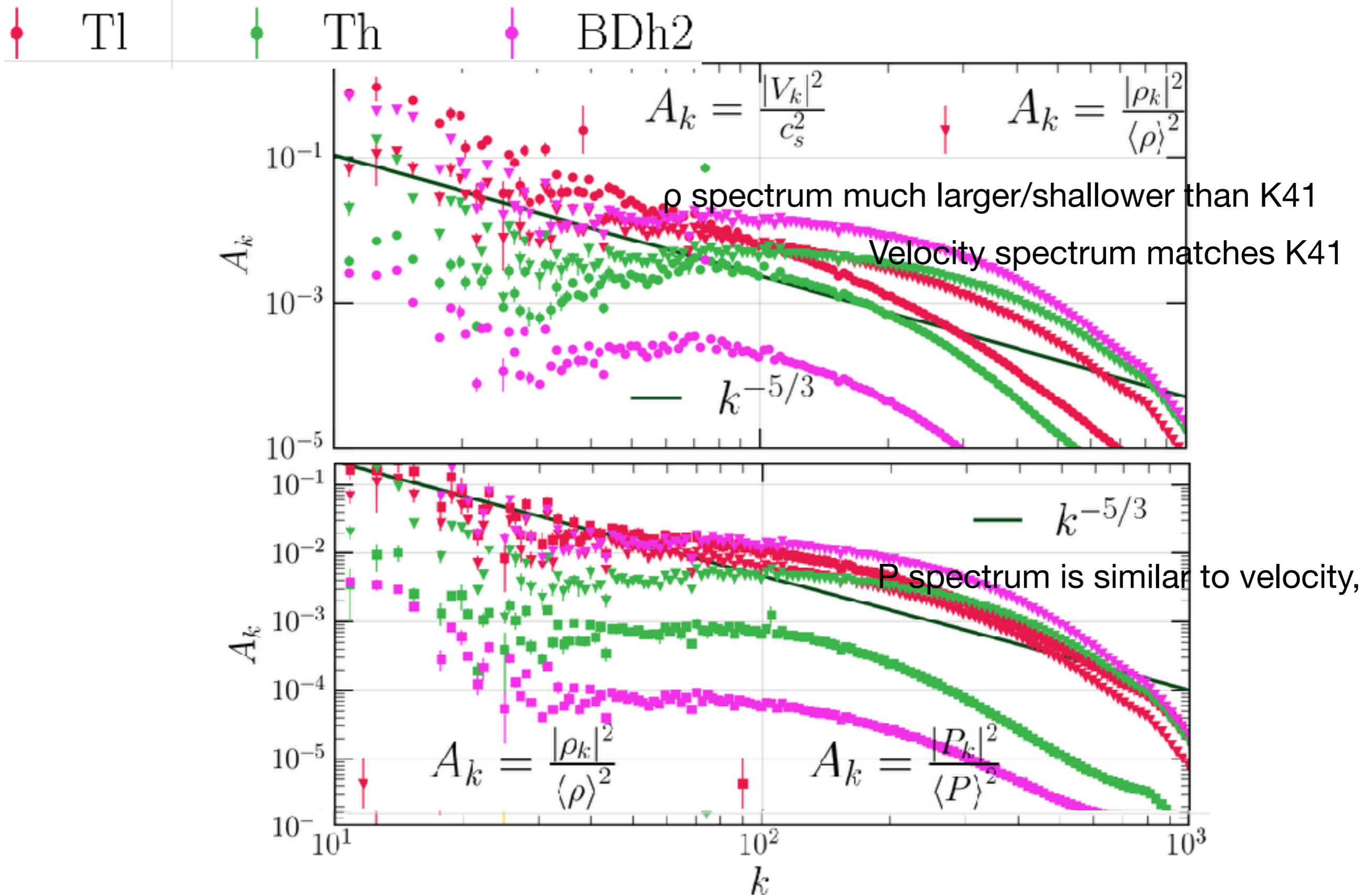
M, T PDF for $f_{\text{turb}} \sim 0.1$



ρ/P fluctuations vs M



Power spectra w htg/clg



Conclusions

- cooling+heating can lead to large density fluctuations in CCs; TI must be accounted for in converting density fluctuations to v_{rms}
- $M < 1$ implies that $f_{turb} \ll 1$ if cold gas is to condense out of the ICM
- $f_{turb} \sim 0.1$ to match Hitomi v_{LOS} in our setup
- Density power spectrum larger/shallower with htg/clg
- Role of **background stratification** in ρ - v_{rms} relation? t_{cool}/t_{ff} ?
- Synergies w high resolution tSZ, X-ray and direct velocity measurements (Hitomi successor)
- MHD, plasma effects!

Thank You!