

Angular momentum transport in accretion

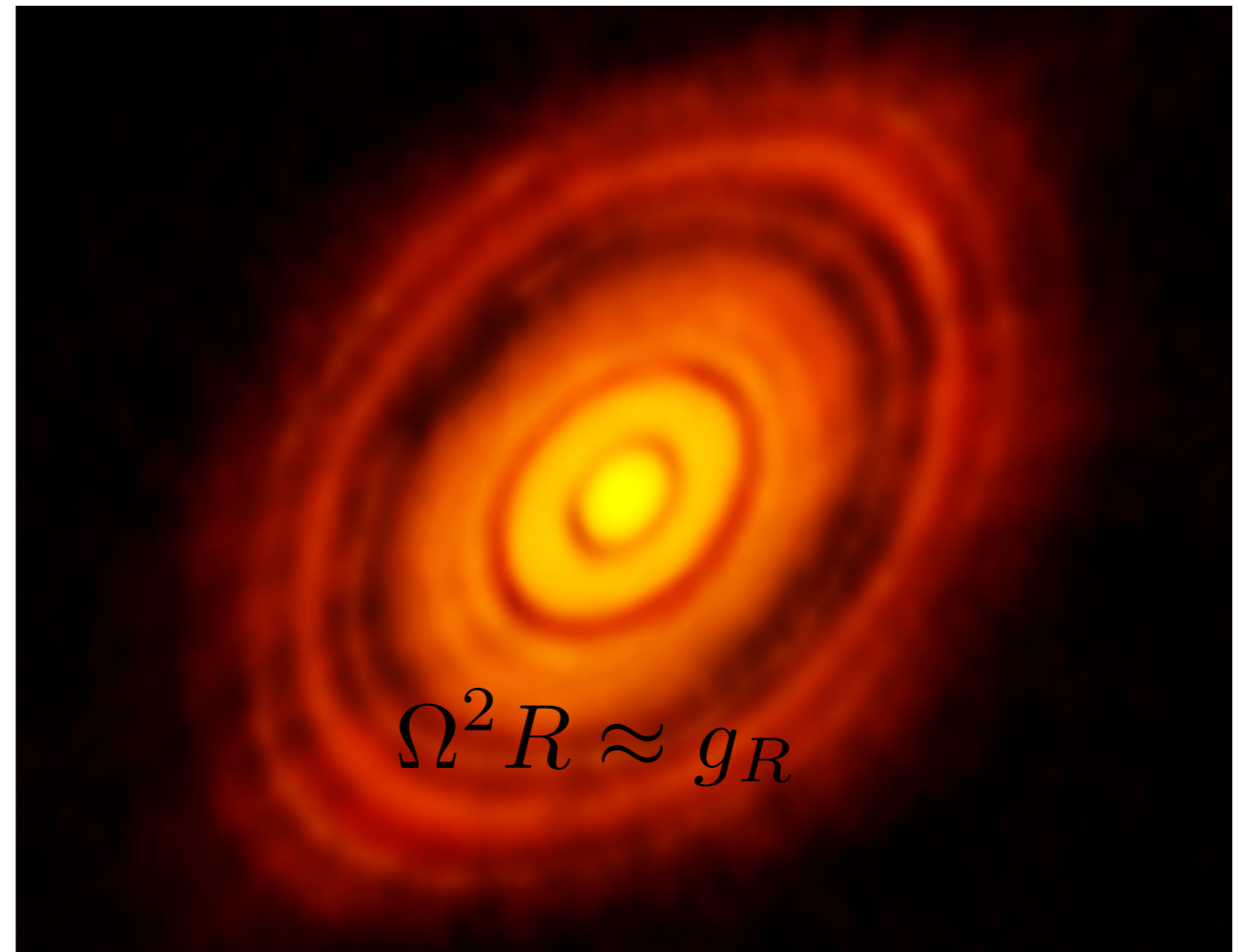
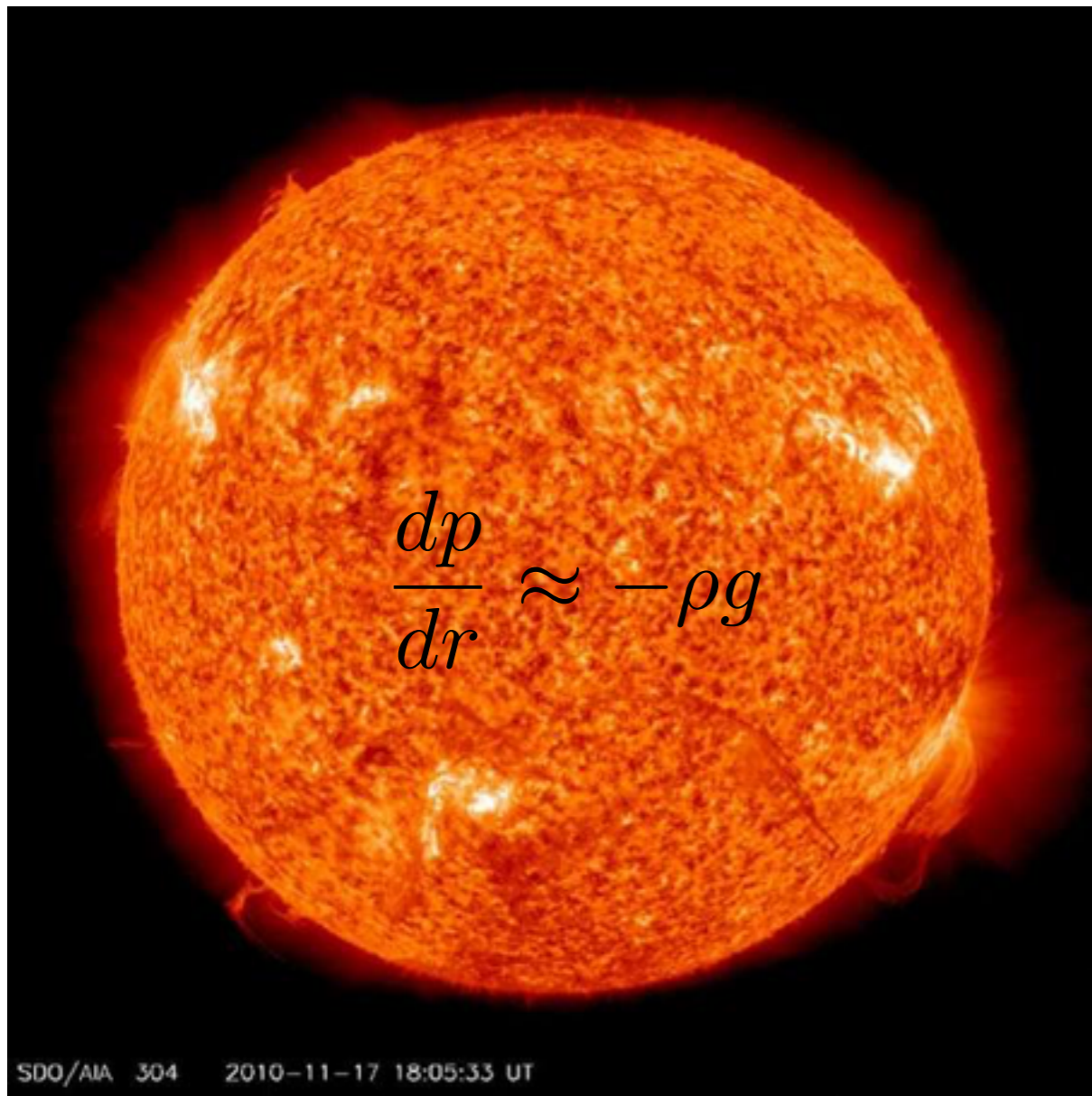
Prateek Sharma, IISc

Accretion

Two major classes of objects in astrophysics:

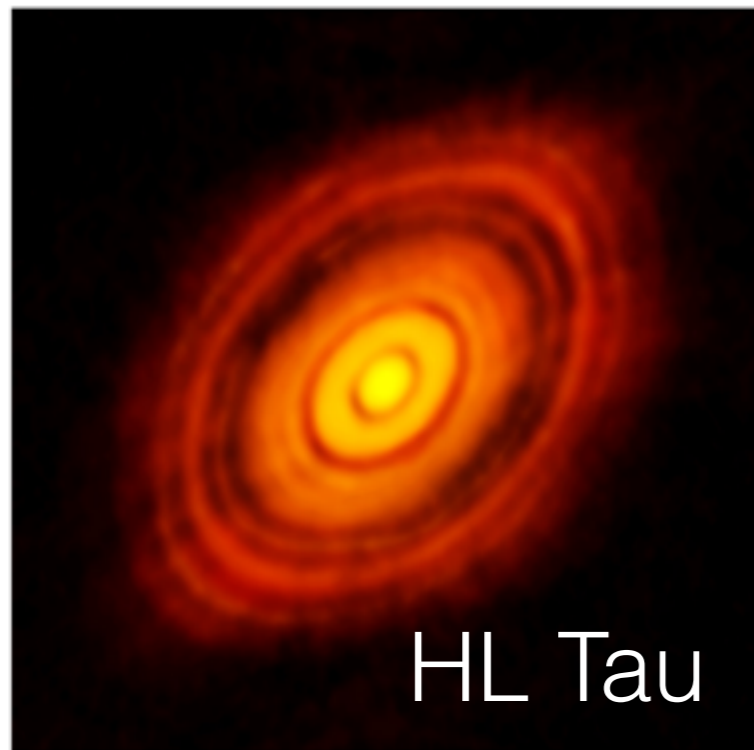
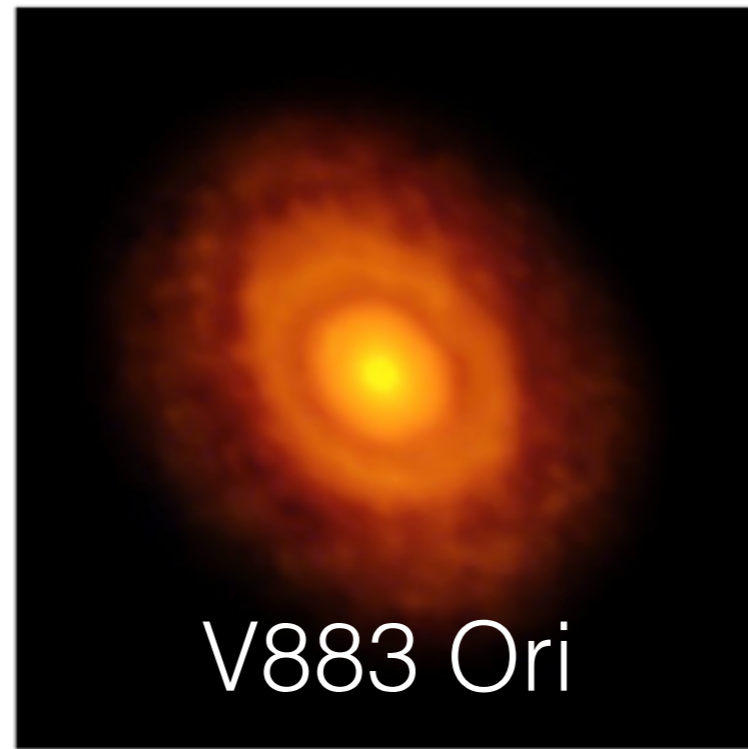
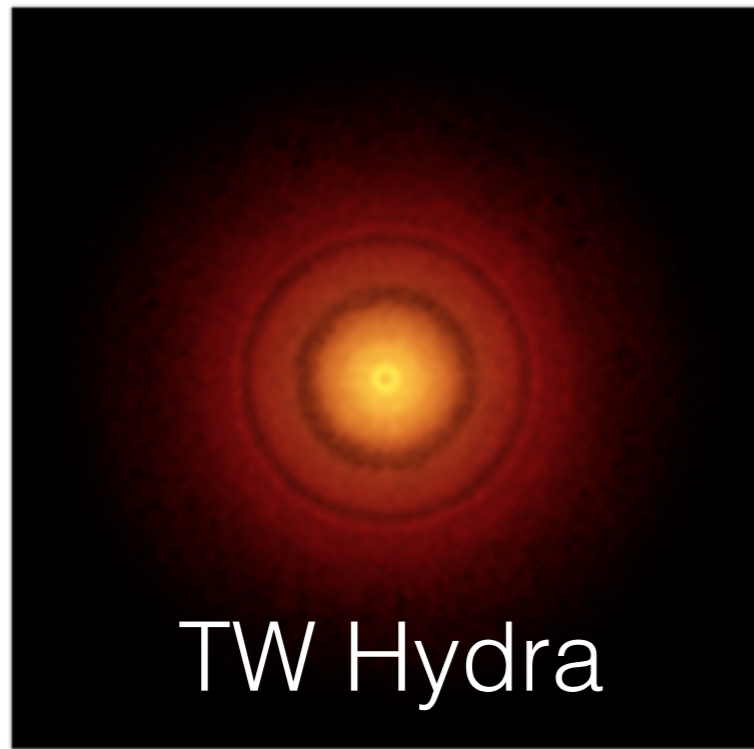
pressure supported, e.g. stars

rotation supported, e.g. disks



formed because collapsing matter lost energy but not angular momentum

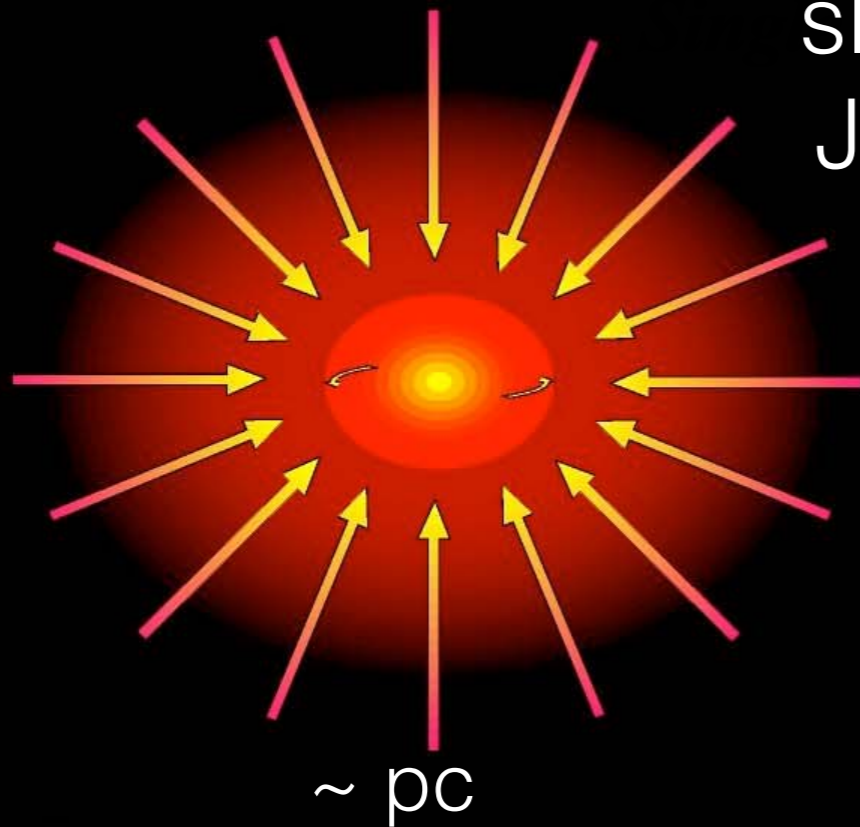
ALMA disks



sub-mm emission from dust in protoplanetary disks ~ 100 AU

Scenario for star- and planet formation

slide courtesy:
Jack Lissaeur

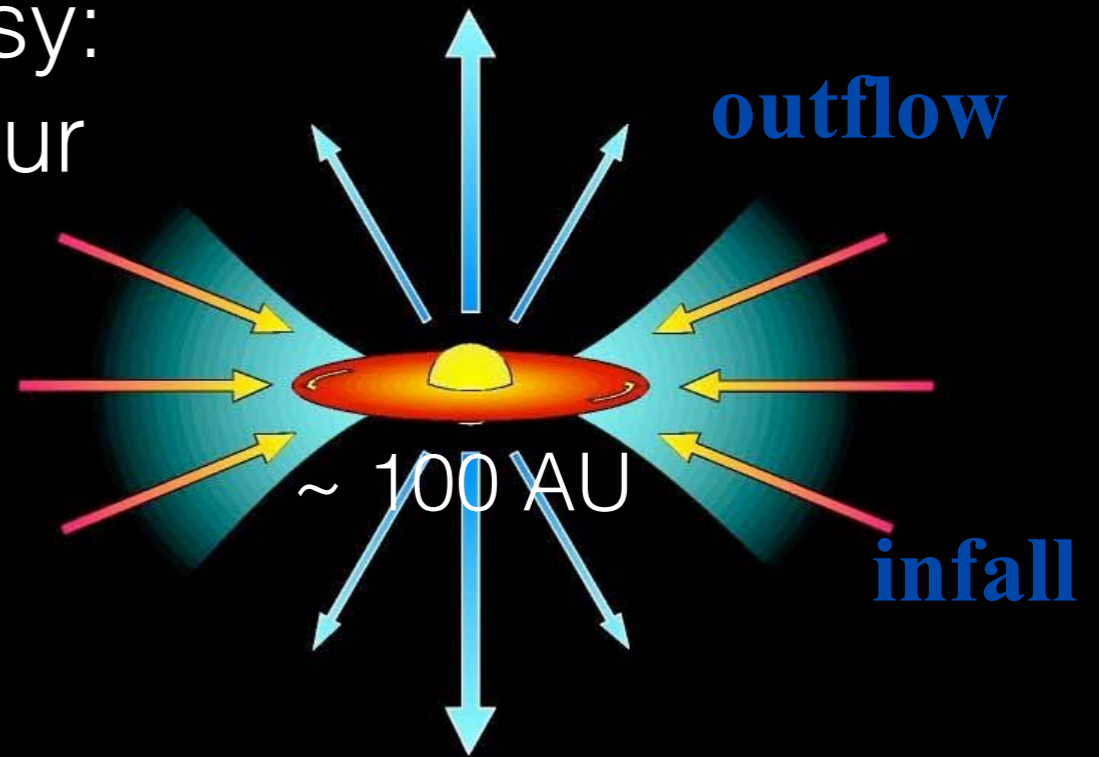


~ pc

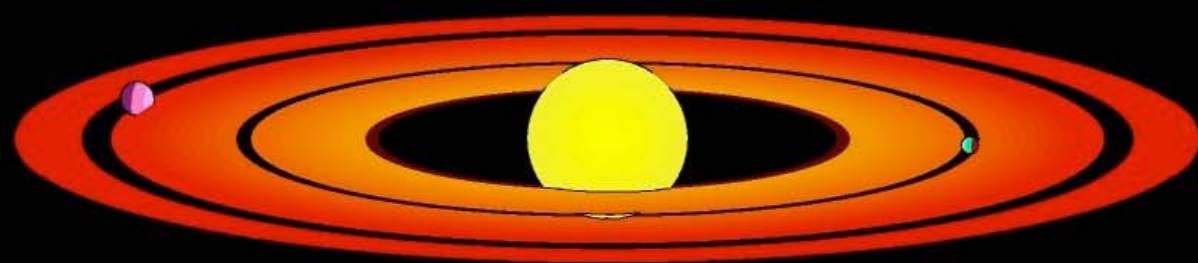
Cloud collapse



**Factor 1000
smaller**

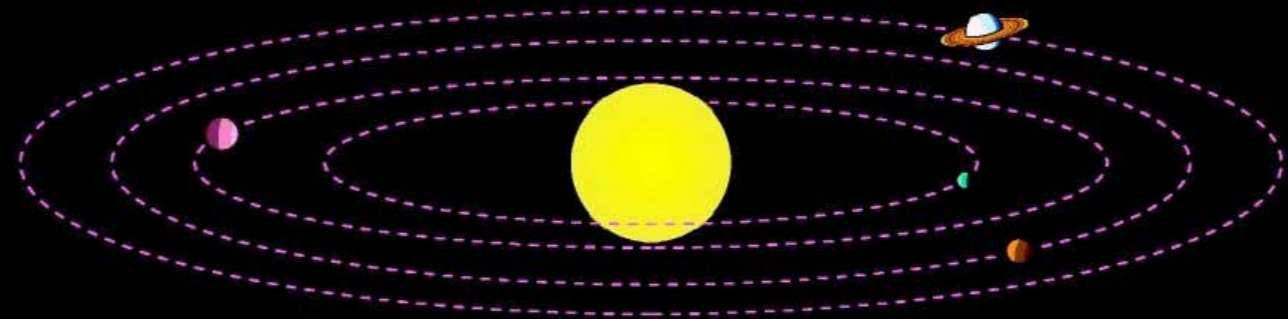


Protostar with disk $t=10^5$ yr



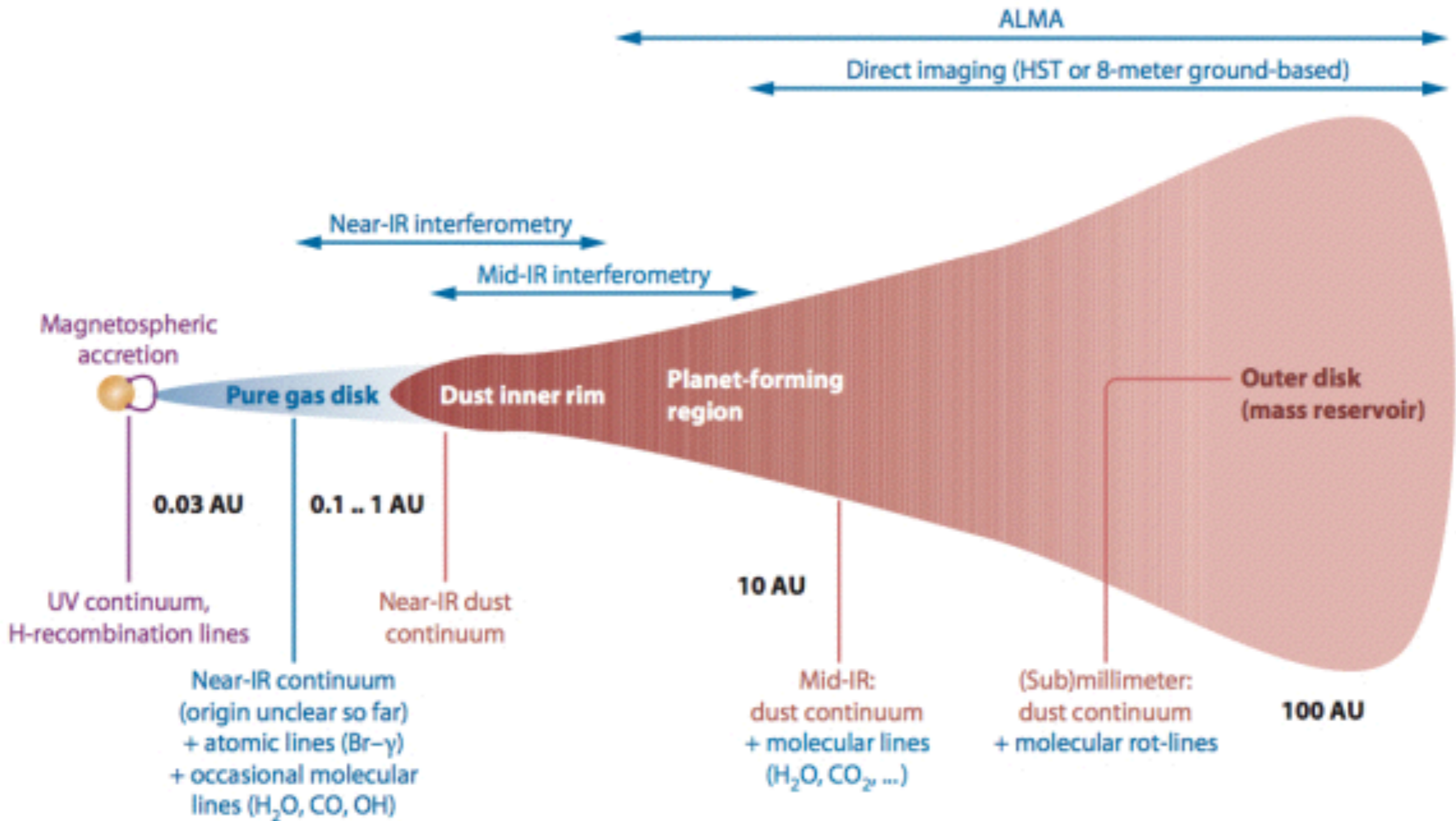
Formation planets

$t=10^6-10^7$ yr



Planetary system $t > 10^8$ yr

Protoplanetary disk

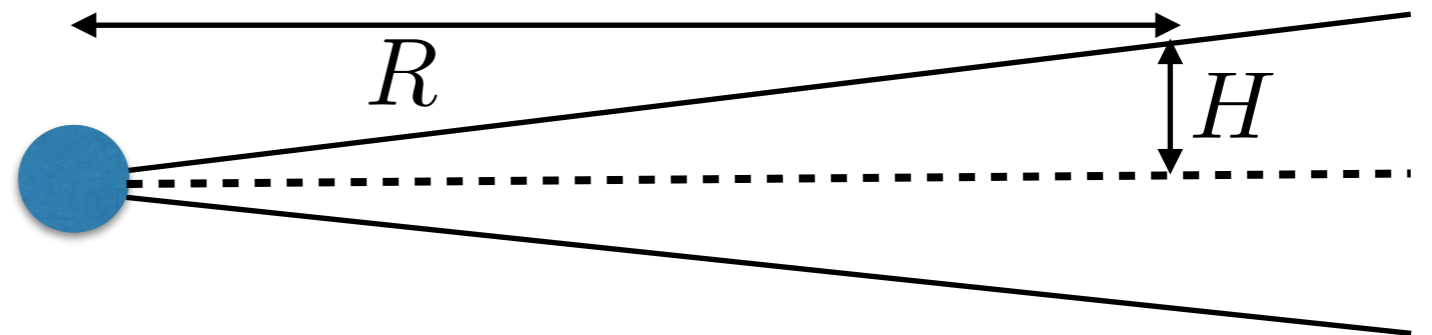


Equilibrium structure

gravity dominated by central object

$$\frac{GM}{R^2} \approx \Omega^2 R \quad \text{radial force balance} \Rightarrow \text{Keplerian flow}$$

$$\Omega_K = \sqrt{\frac{GM}{R^3}} \propto R^{-3/2}$$

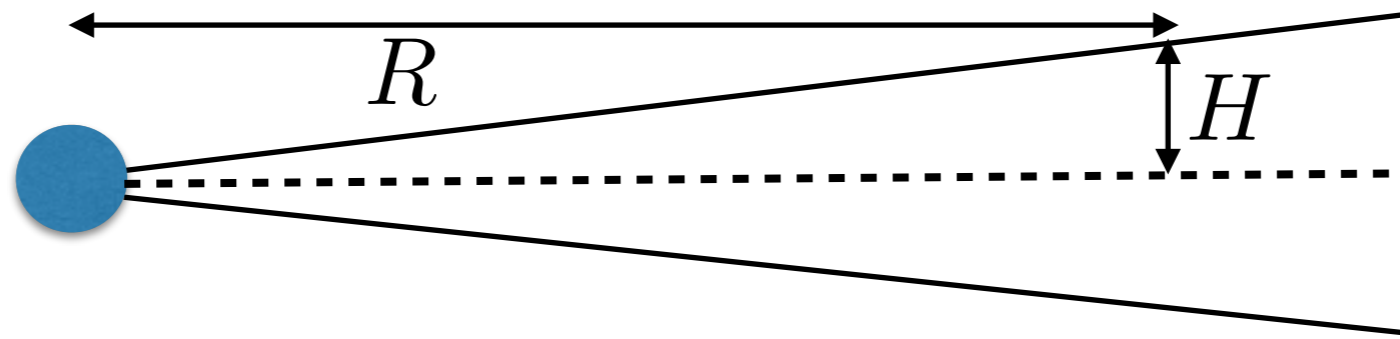


$$\frac{1}{\rho} \frac{dp}{dz} \approx -\frac{GMz}{R^3} = -\Omega_K^2 z \quad \text{vertical HSE} \Rightarrow H \approx c_s / \Omega_K$$

$$H/R \approx c_s / v_K \ll 1$$

dynamical equilibrium & **slow** radial accretion of matter
radial & vertical dynamics decoupled \Rightarrow 1-D solutions vs.
R of z-integrated quantities

Energy considerations

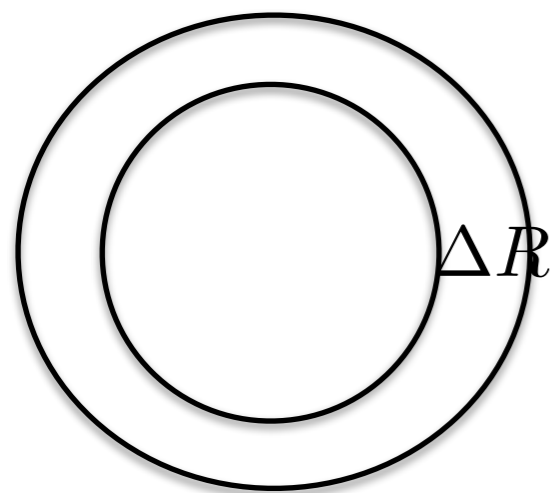


total energy per mass for an orbiting particle = $-\frac{GM}{2R}$

KE in rotation = $\frac{GM}{2R}$ gravitational PE = $-\frac{GM}{R}$

$\frac{GM}{2R}$ energy per accreted mass available as heat & light

accretion power generated from R to $R+\Delta R$:



$$GM\dot{M}\Delta\left(\frac{-1}{2R}\right) = \frac{GM\dot{M}}{2R^2}\Delta R$$

total power released: $\frac{GM\dot{M}}{2R_{\text{in}}}$

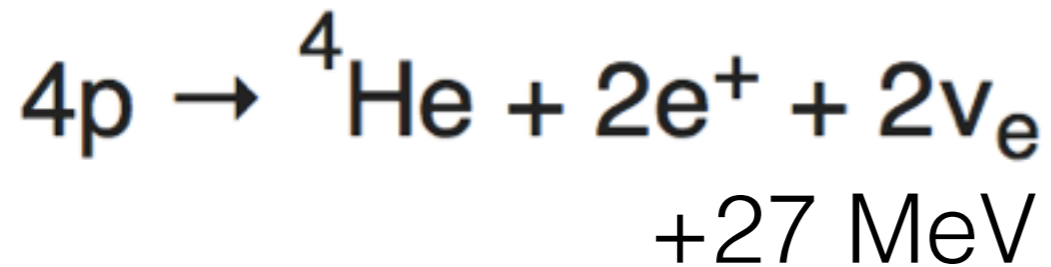
Accretion power

$$\frac{GM\dot{M}}{2R_{\text{in}}} = \eta\dot{M}c^2$$

$$\eta = \frac{GM}{2R_{\text{in}}c^2} = \frac{v_{\text{esc}}^2}{4c^2} = \frac{R_{\text{Sch}}}{4R_{\text{in}}}$$

higher efficiency for compact accretor

accretion efficiency

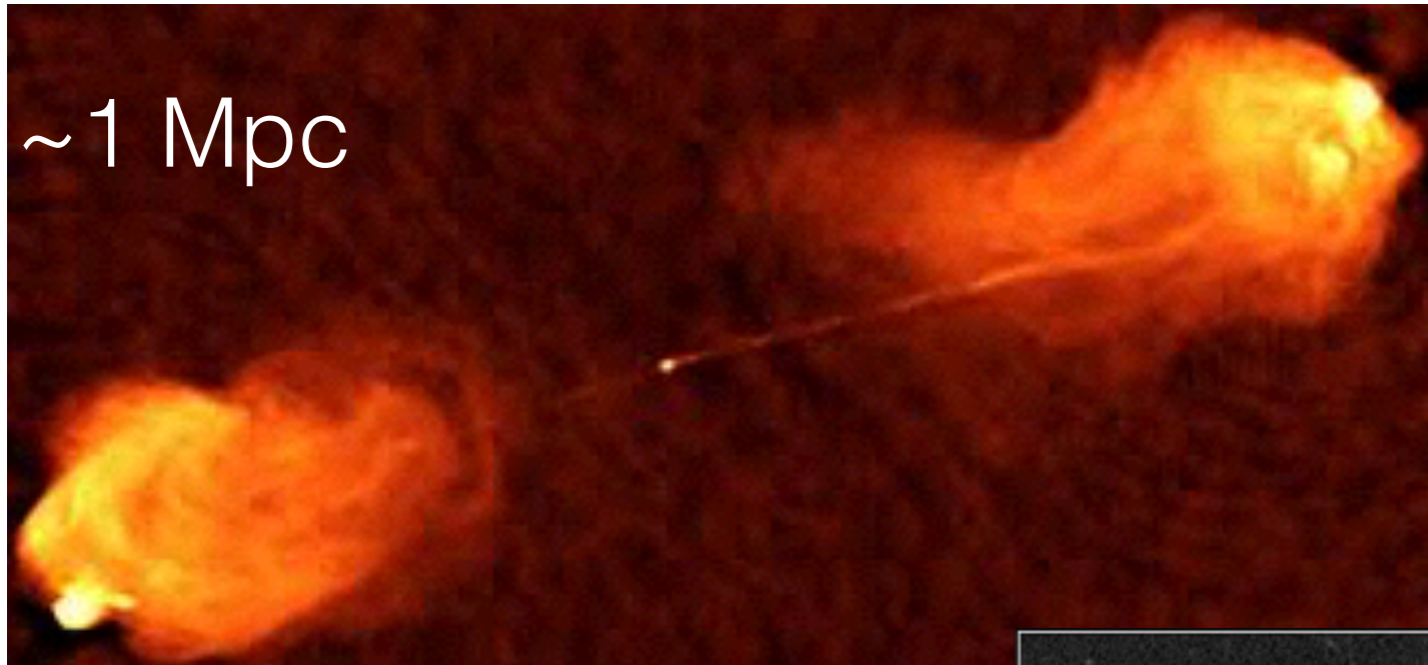


$$\eta_{\text{nuc}} \approx 27 \text{ MeV} / 4 \text{ GeV} = 0.007$$

no wonder most powerful sources due to accretion!

	$R_{\text{in}}(\text{km})$	$M(M_{\odot})$	$R_{\text{in}}/R_{\text{Sch}}$	η
protostar	10^6	1	3×10^5	10^{-6}
WD	10^4	1	3000	10^{-4}
NS	10	1	3	0.1
BH	$\frac{9M}{M_{\odot}}$	10, 10^6 -9	3	0.1

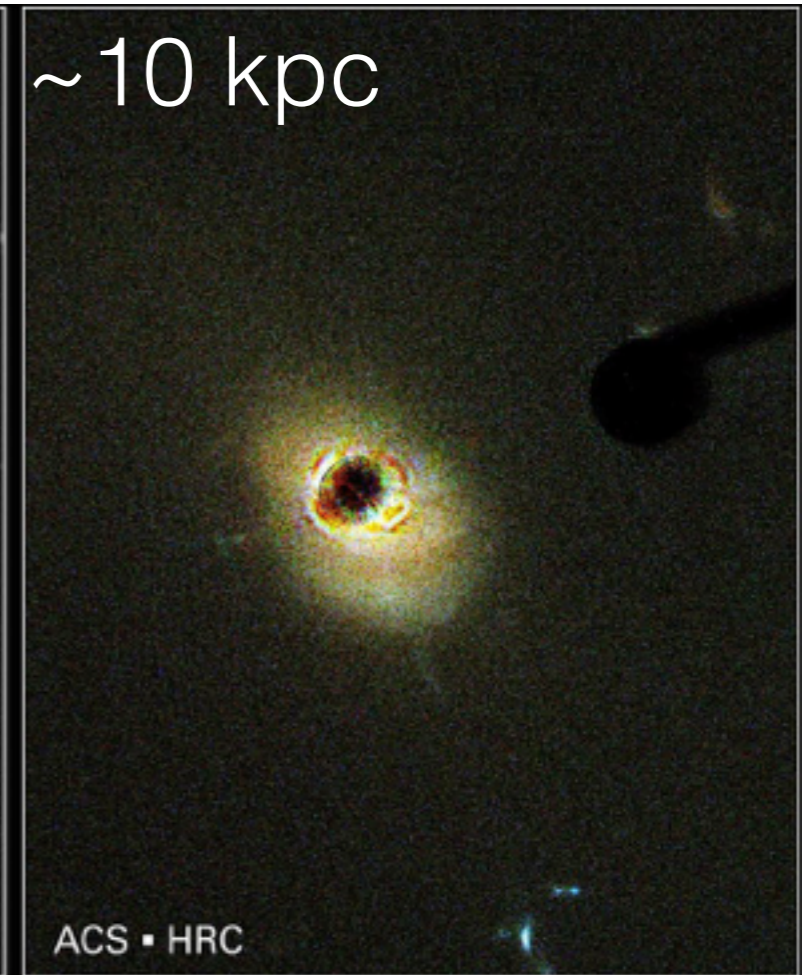
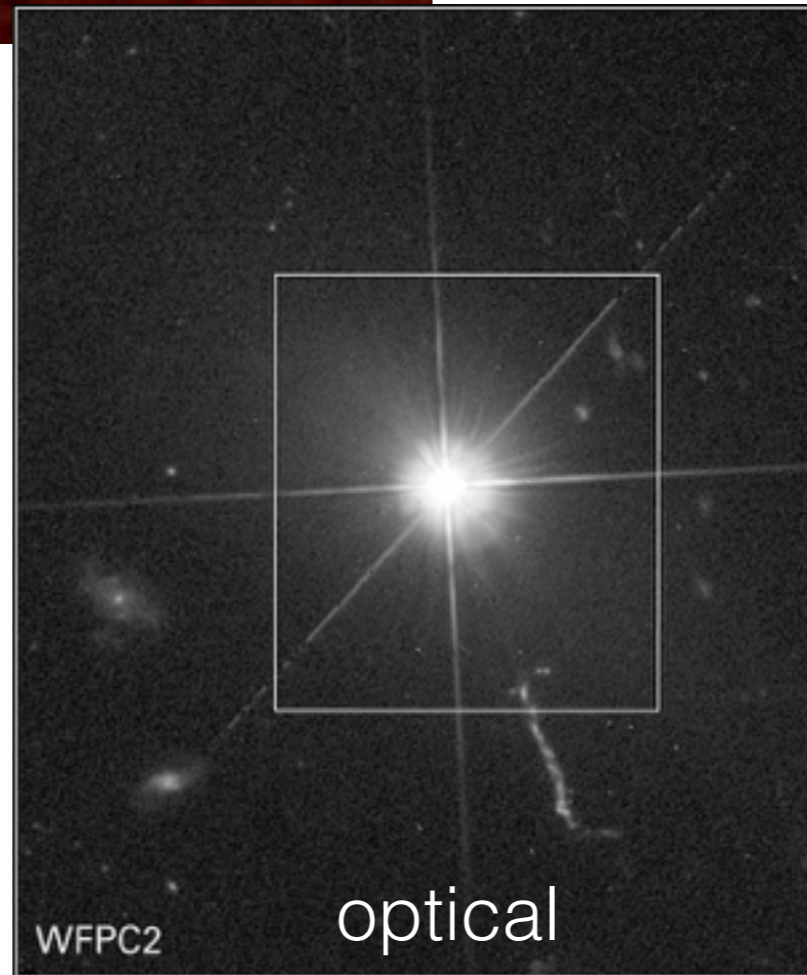
Manifestation



~1 Mpc

quasar 3C-273 in a galaxy
radiative power

radio galaxy, Cyg A
mechanical power
synchrotron radiation
from relativistic e-s



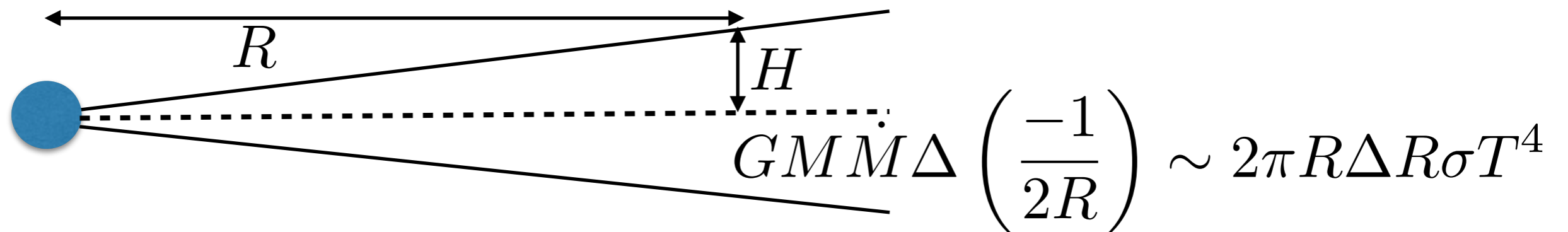
Flow Temperature

Two estimates:

virial applicable for non-radiative disks

$$k_B T_{\text{virial}} \equiv GMm_p/3R \sim 0.1m_p c^2 \sim 10^{12} \text{K}$$

geometrically thin, optically thick disk, a local BB disk



$T \sim 1$ keV for maximally accreting 10 solar mass BH, emits X-rays

Angular momentum transport

Key Q: how does matter lose angular momentum and fall in?
“molecular viscosity” negligible, need turbulent transport, **but how?**

Linear response: $w^2 = \frac{k_z^2}{k^2} \kappa^2$, $\kappa^2 = \frac{2\Omega}{R} \frac{dl}{dR}$

$$l = \sqrt{GMR}$$

stable inertial waves
or epicycles oscillations

a perturbed fluid element preserves angular momentum

& **hydro Keplerian disks are Rayleigh stable**

(to local axisymmetric modes)

Hydro Keplerian flow does not naturally break into turbulence

Linear MHD instability

[Balbus & Hawley 1991]

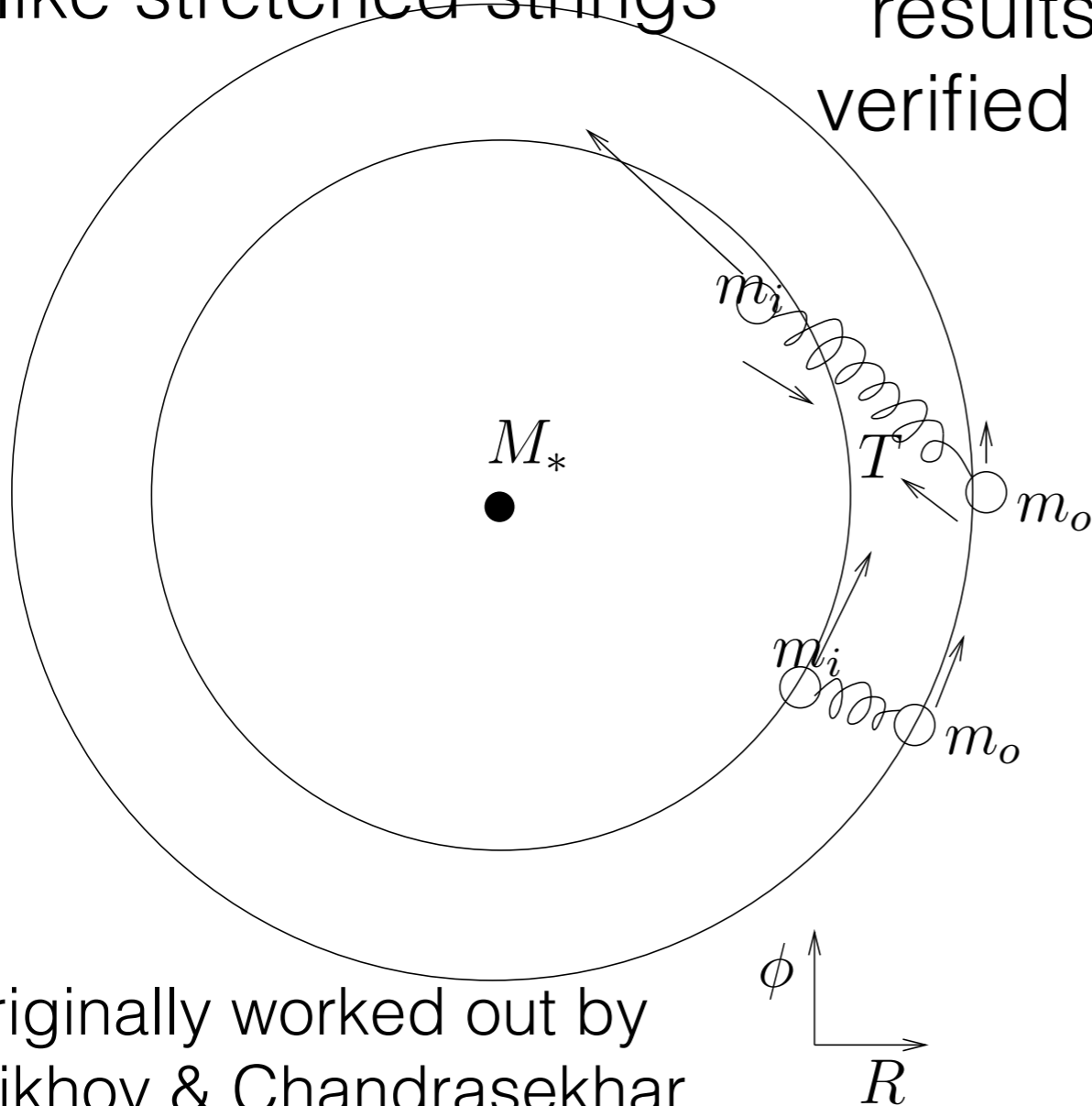
local axisymmetric MHD \Rightarrow MRI

magnetorotational instability

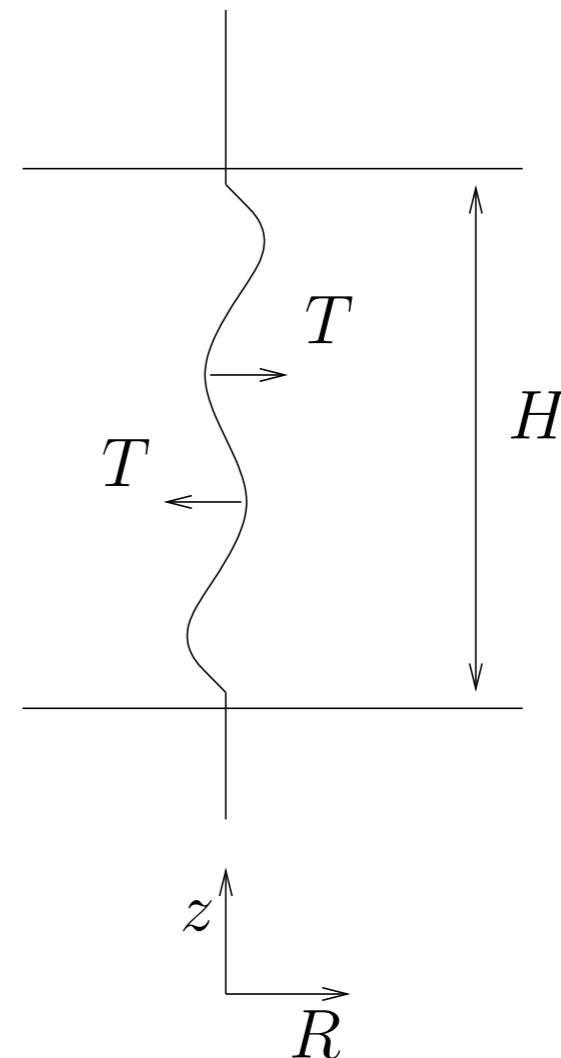
works for ionized flows

results in **MHD turbulence & transport**
verified by local & global MHD sims [DNS]

magnetic fields behave
like stretched strings



originally worked out by
Velikhov & Chandrasekhar

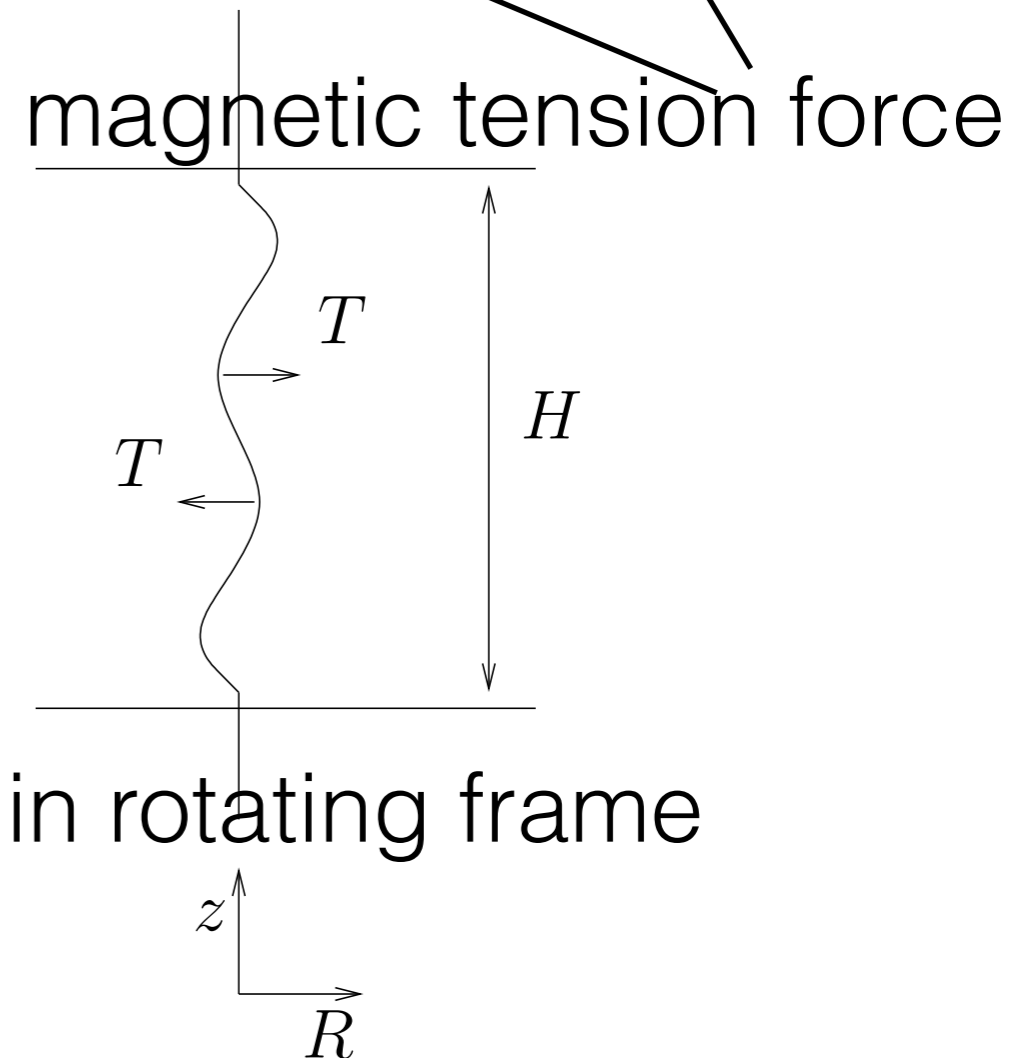
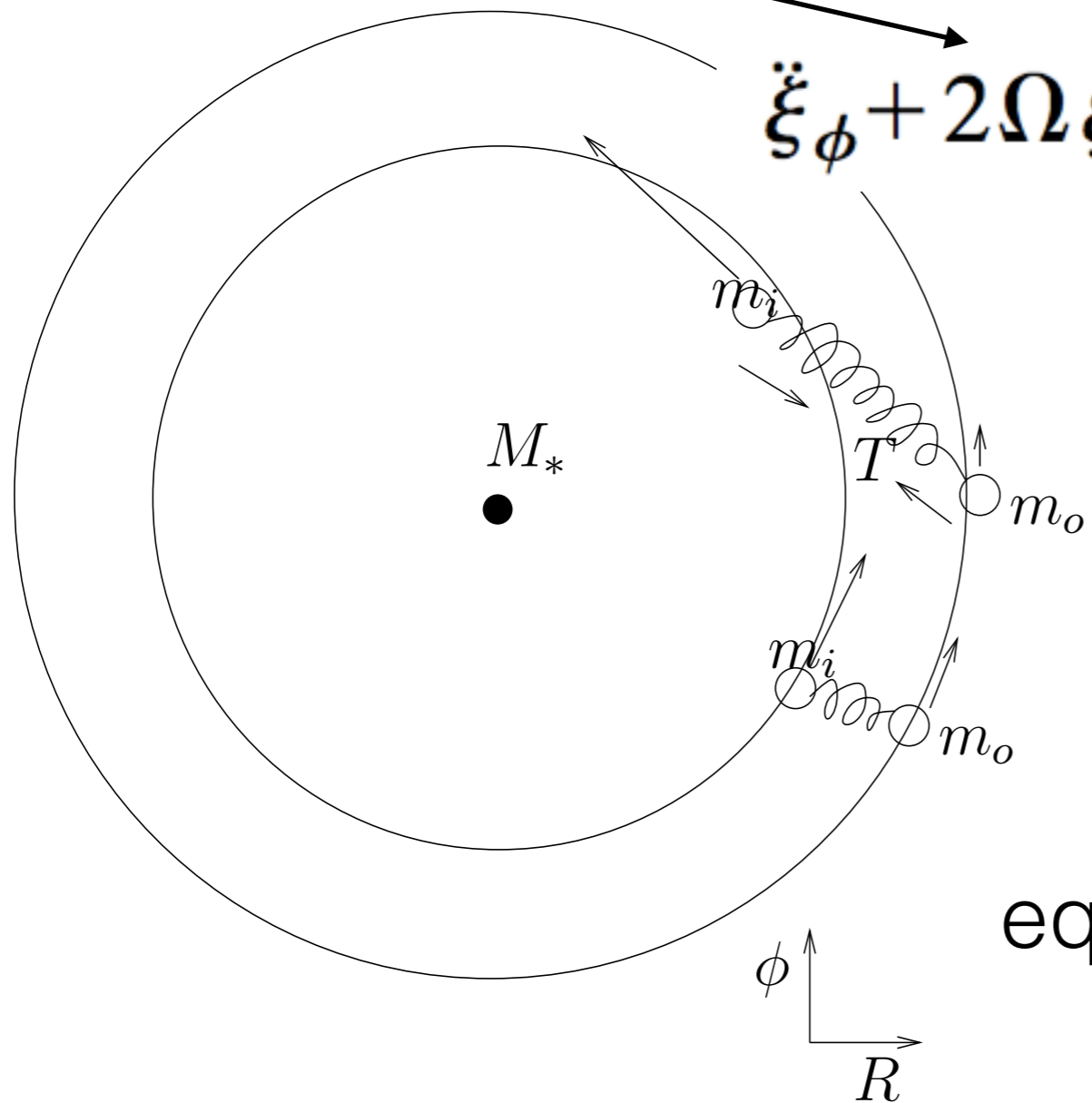


Linear MHD instability

[Balbus & Hawley 1991]

$$\ddot{\xi}_R - 2\Omega \dot{\xi}_\phi = - \left(\frac{d\Omega^2}{d \ln R} + (\mathbf{k} \cdot \mathbf{u}_A)^2 \right) \xi_R,$$

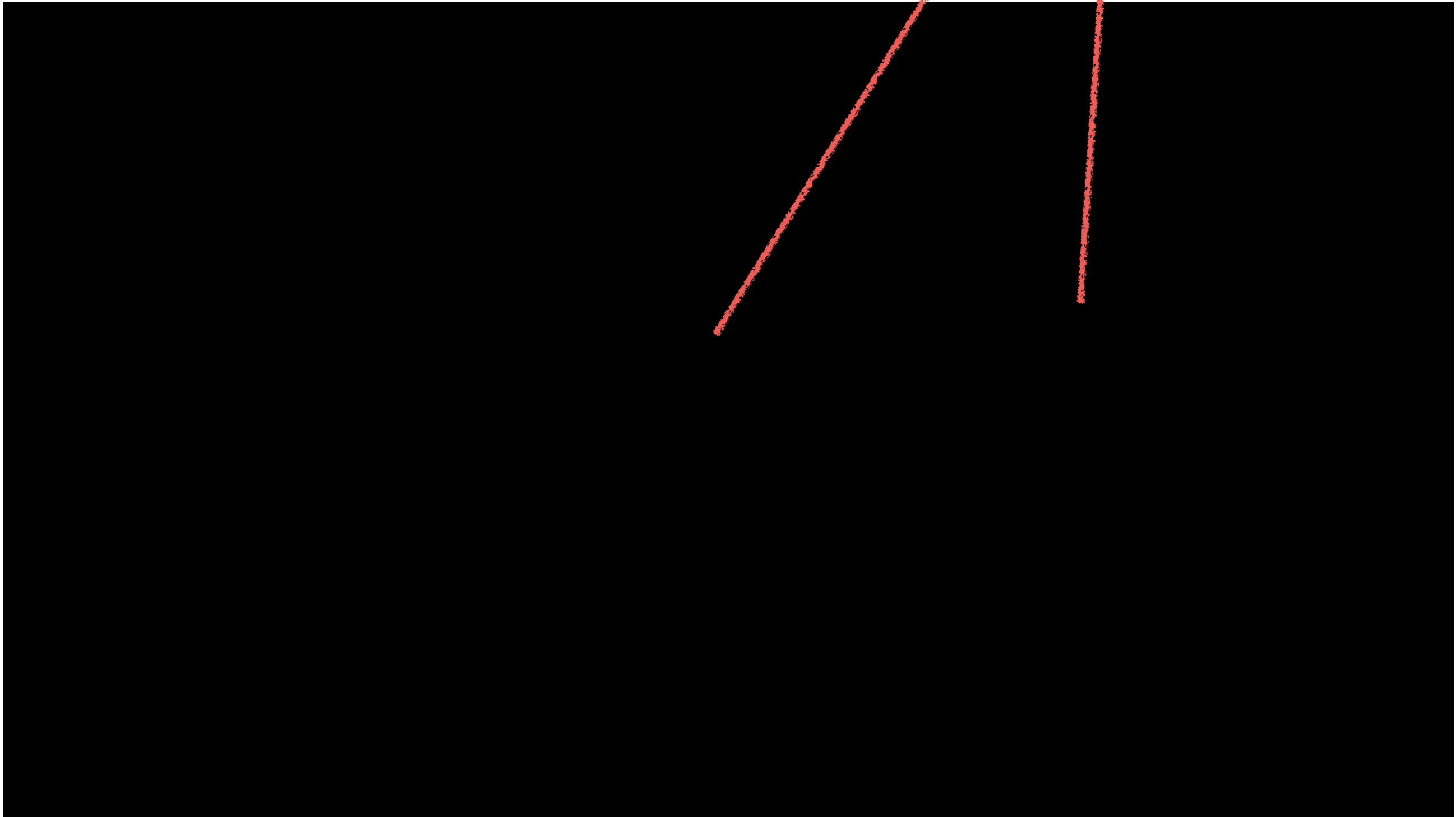
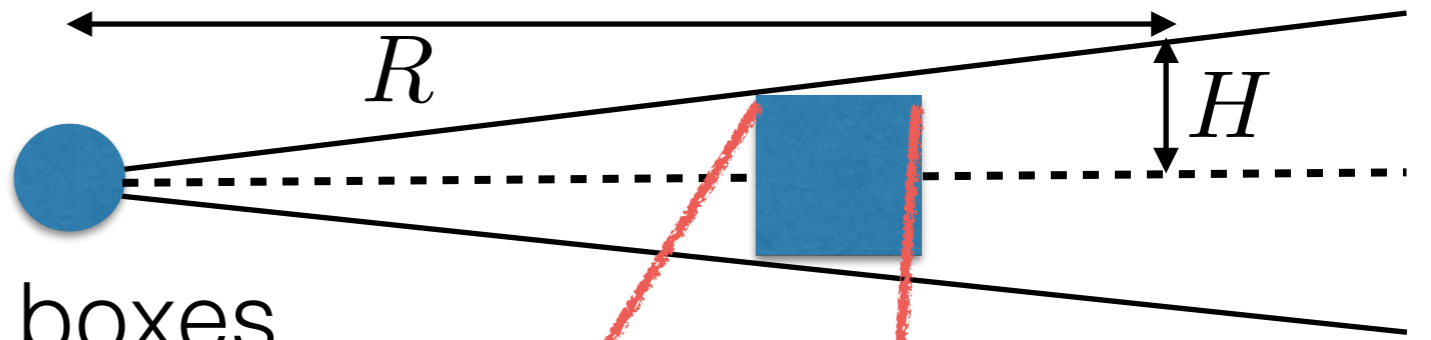
$$\ddot{\xi}_\phi + 2\Omega \dot{\xi}_R = - (\mathbf{k} \cdot \mathbf{u}_A)^2 \xi_\phi.$$



equations in rotating frame

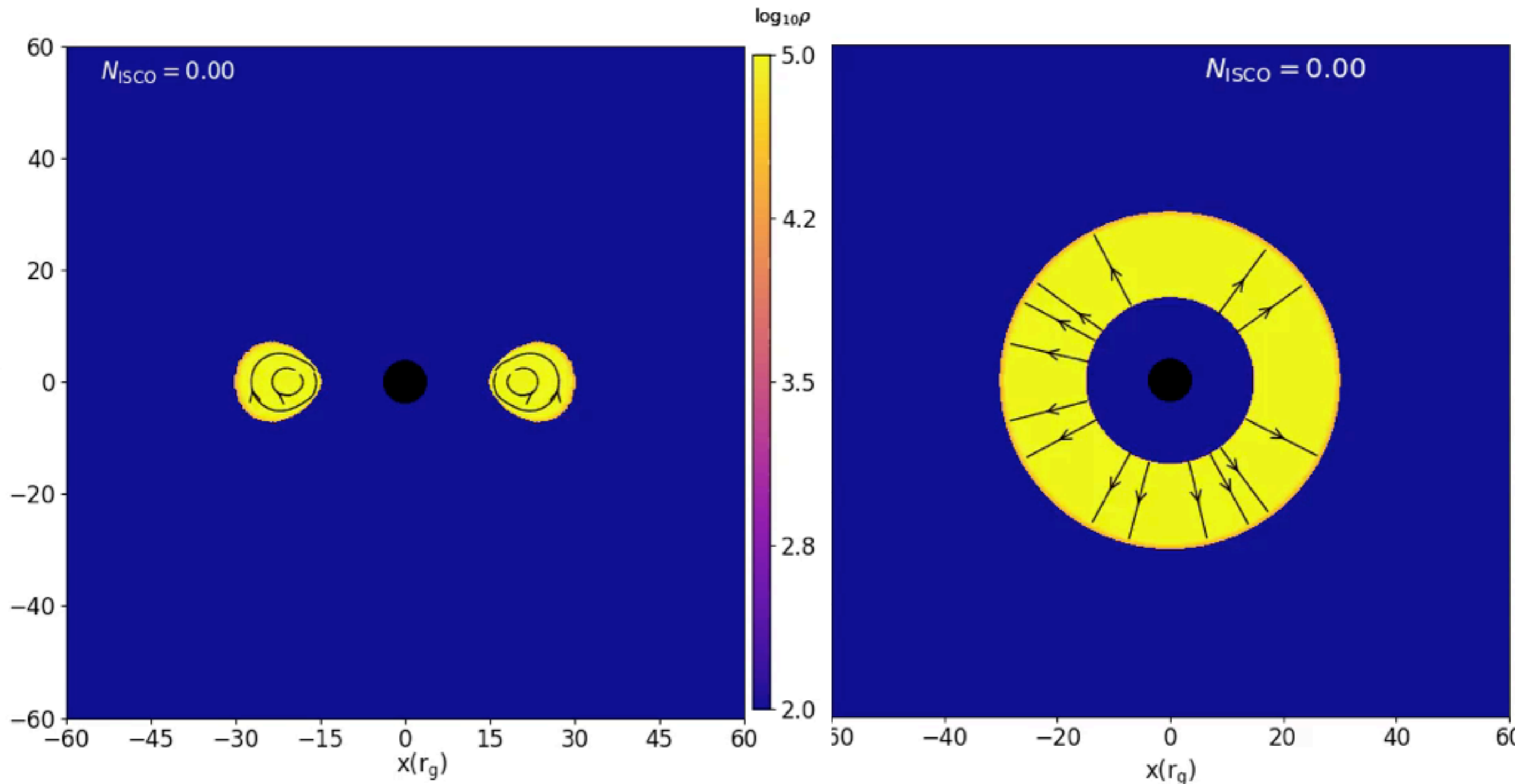
MHD sims

local shearing boxes

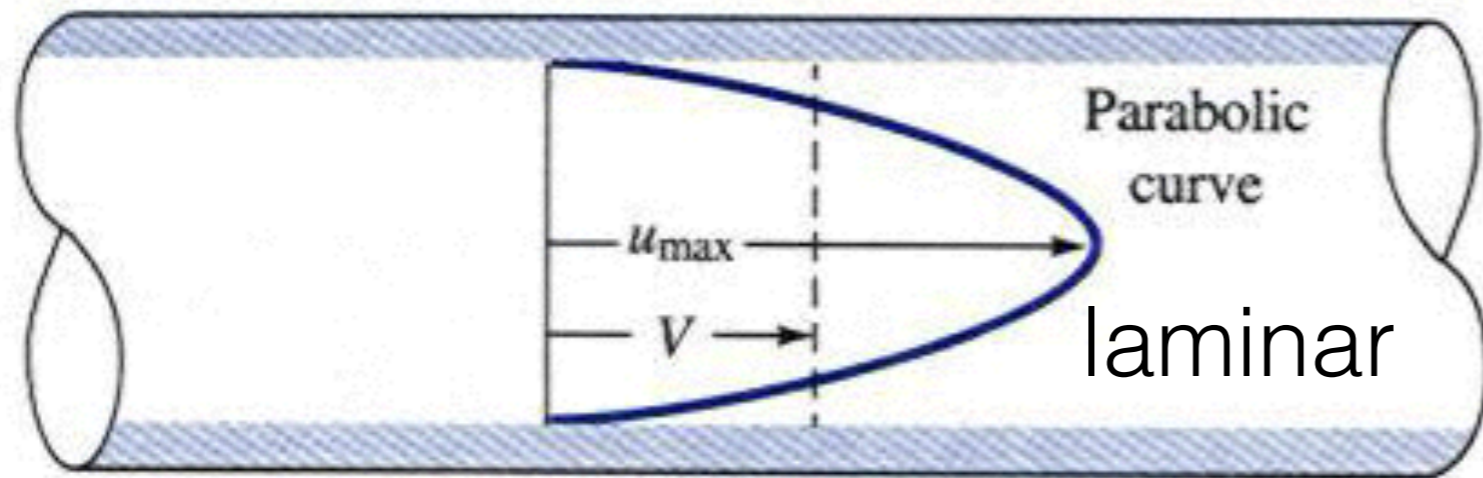


Global simulations

[movies by Prasun Dhang]

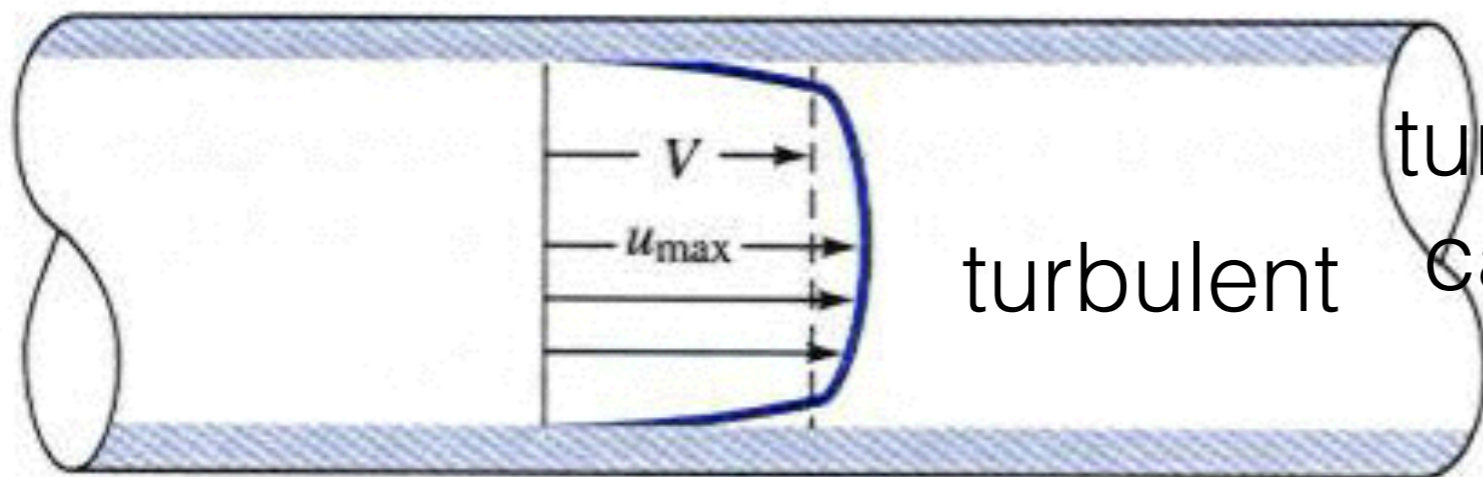


Turbulent transport



(a)

Pipe flow:
stream-wise turbulent
momentum transport
 \Rightarrow flattening of velocity
profile



(b)

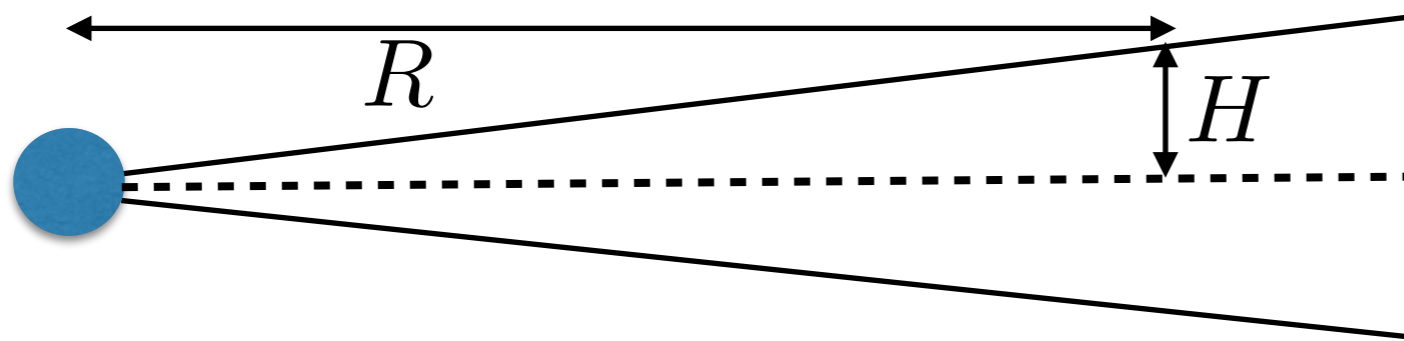
turbulent correlations $\langle U_R U_Z \rangle$
cause momentum transport
from center to wall

Turbulent Reynolds stress modifies mean flow

Angular momentum eq.

inviscid, ϕ -averaged angular momentum equation

$$\frac{\partial}{\partial t}(\rho \langle u_\phi \rangle R) + \frac{1}{R} \frac{\partial}{\partial R} \left[R^2 \left(\langle \rho u_\phi u_R - \frac{B_\phi B_R}{4\pi} \rangle \right) \right] + \frac{\partial}{\partial z} \left(R \langle \rho u_\phi u_z - \frac{B_\phi B_z}{4\pi} \rangle \right) = 0$$



$$\Sigma \equiv \int_{-H/2}^{H/2} \rho dz$$

$$\dot{M} \equiv 2\pi R \Sigma \langle u_R \rangle$$

in SS, integrating over z , angular momentum flow rate is const.

$$R^2 \left[\Sigma \langle u_R \rangle \langle u_\phi \rangle + \langle \Sigma u_{R1} u_{\phi 1} - H \frac{B_R B_\phi}{4\pi} \rangle \right]$$

Reynolds & Maxwell stress

$$W_{R\phi} = \frac{\dot{M} l}{2\pi R^2 H} \left(1 - \frac{l_\star}{l} \right)$$

non-zero turbulent stress
needed for mass accretion!

Angular momentum

specific angular momentum $l = \sqrt{GM R} \propto R^{1/2}$

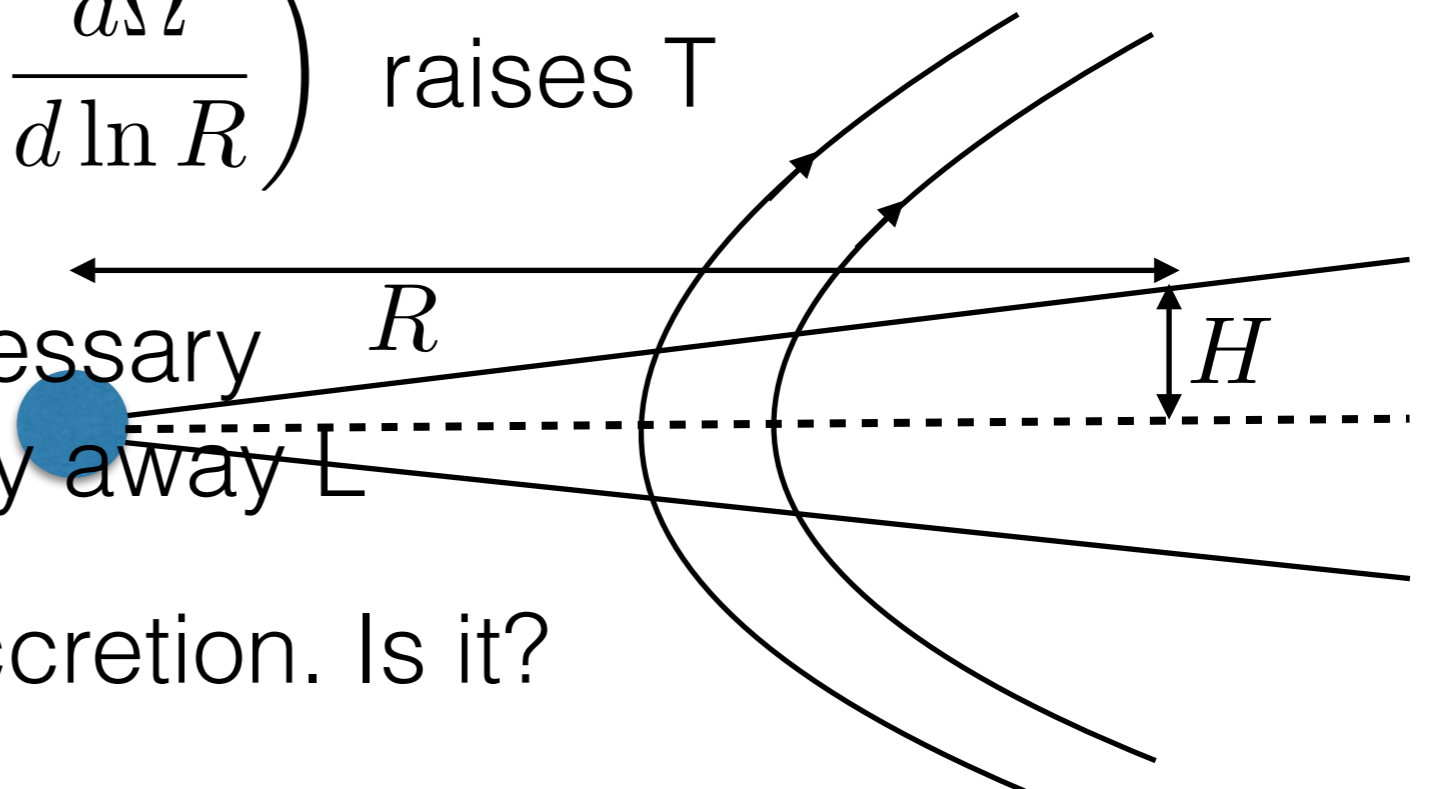
increases with R ; since friction is an internal force only small mass can carry away most angular momentum

$$\dot{M}_{\text{out}} \sim \dot{M}_{\text{in}} \left(\frac{R_{\text{in}}}{R_{\text{out}}} \right)^{1/2}$$

viscous heating $\rho \nu_{\text{turb}} \left(\frac{d\Omega}{d \ln R} \right)^2$ raises T

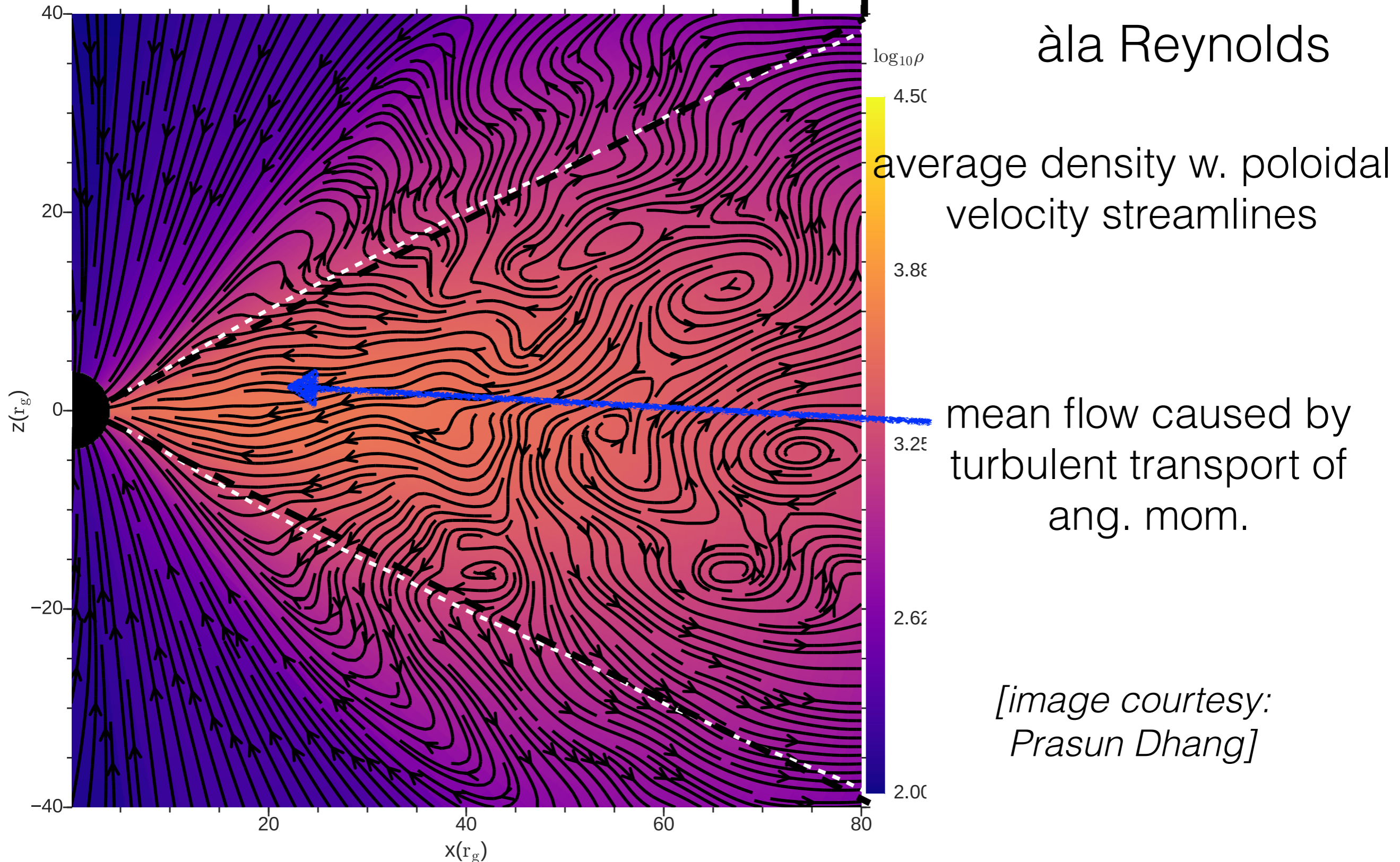
disk turbulence not necessary
magnetized wind can carry away L

MHD looks critical for accretion. Is it?



Mean field approach

à la Reynolds



*[image courtesy:
Prasun Dhang]*

disk temperature

$$T \sim 1 \text{ keV} \left(\frac{M}{10 M_{\odot}} \right)^{1/4} \left(\frac{\dot{M}}{10^{-8} M_{\odot} \text{ yr}^{-1}} \right)^{1/4} \left(\frac{r}{100 \text{ km}} \right)^{-3/4}$$

gives $\sim 100 \text{ K}$ for $1 M_{\text{sun}}$ & $r=1 \text{ AU}$

additional heating due to heating by stellar light

$$\frac{n_i n_e}{n_H} = \frac{(2\pi m_e k_B T)^{3/2}}{h^3} e^{-\chi/k_B T}$$

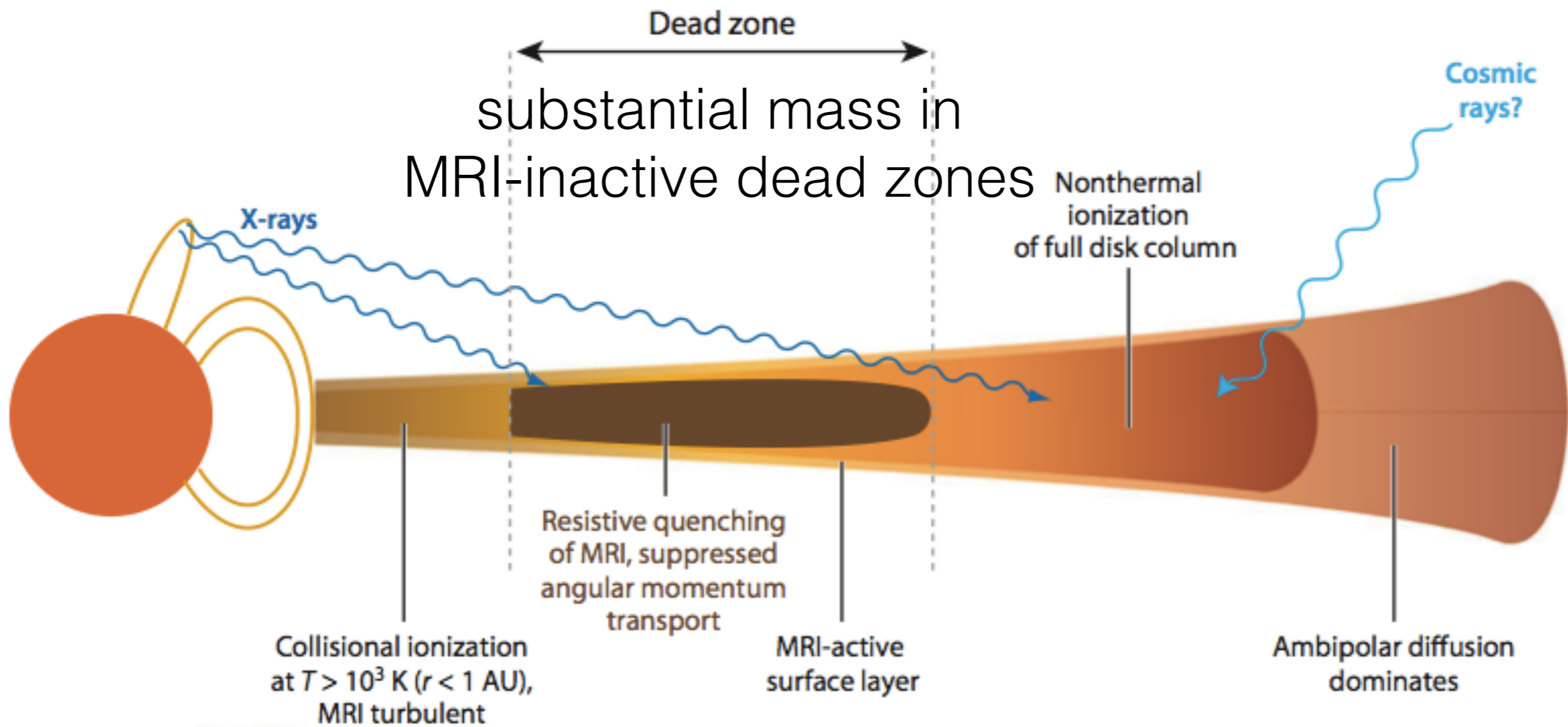
thermal ionization essentially gives $x_e \sim 0!$

proto-planetary disks (PPDs) are completely neutral

non-thermal ionization due to X-ray,

CRs, radioactivity gives $x_e \sim 10^{-13}$

Is hydro turbulence possible?

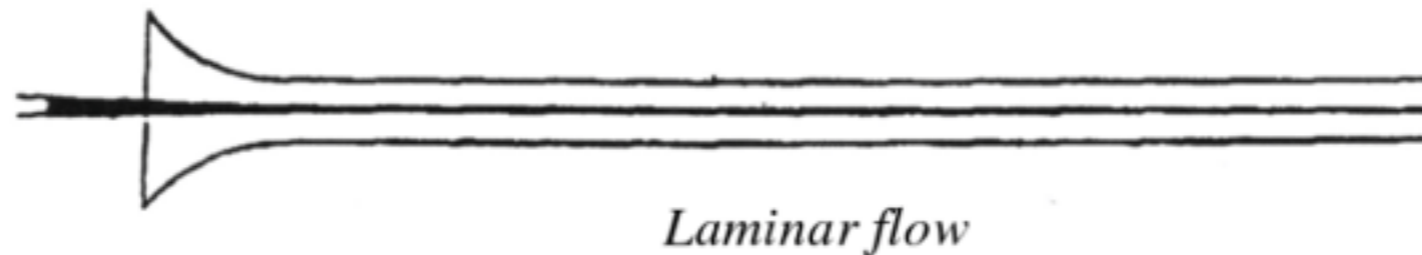


can hydro mechanisms operate?

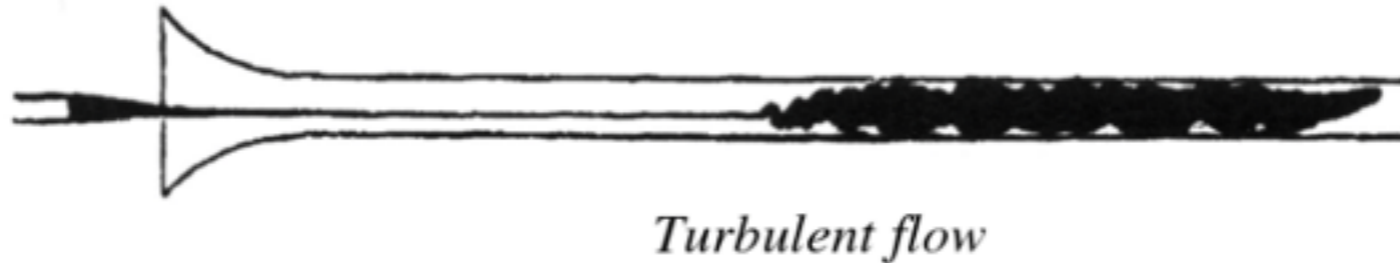
various non-ideal MHD effects: resistivity, ambipolar diffusion, Hall term

What abt hydro transport?

so common for laminar flow to break into turbulence at large Re



pipe flow linearly stable for all Re; transition observed at $Re \sim 2000$
sensitive to roughness, vibrations, etc.



why shouldn't this happen in Keplerian flows where $Re > 10^{10}$?

stabilizing role of (fast) epicyclic oscillations!

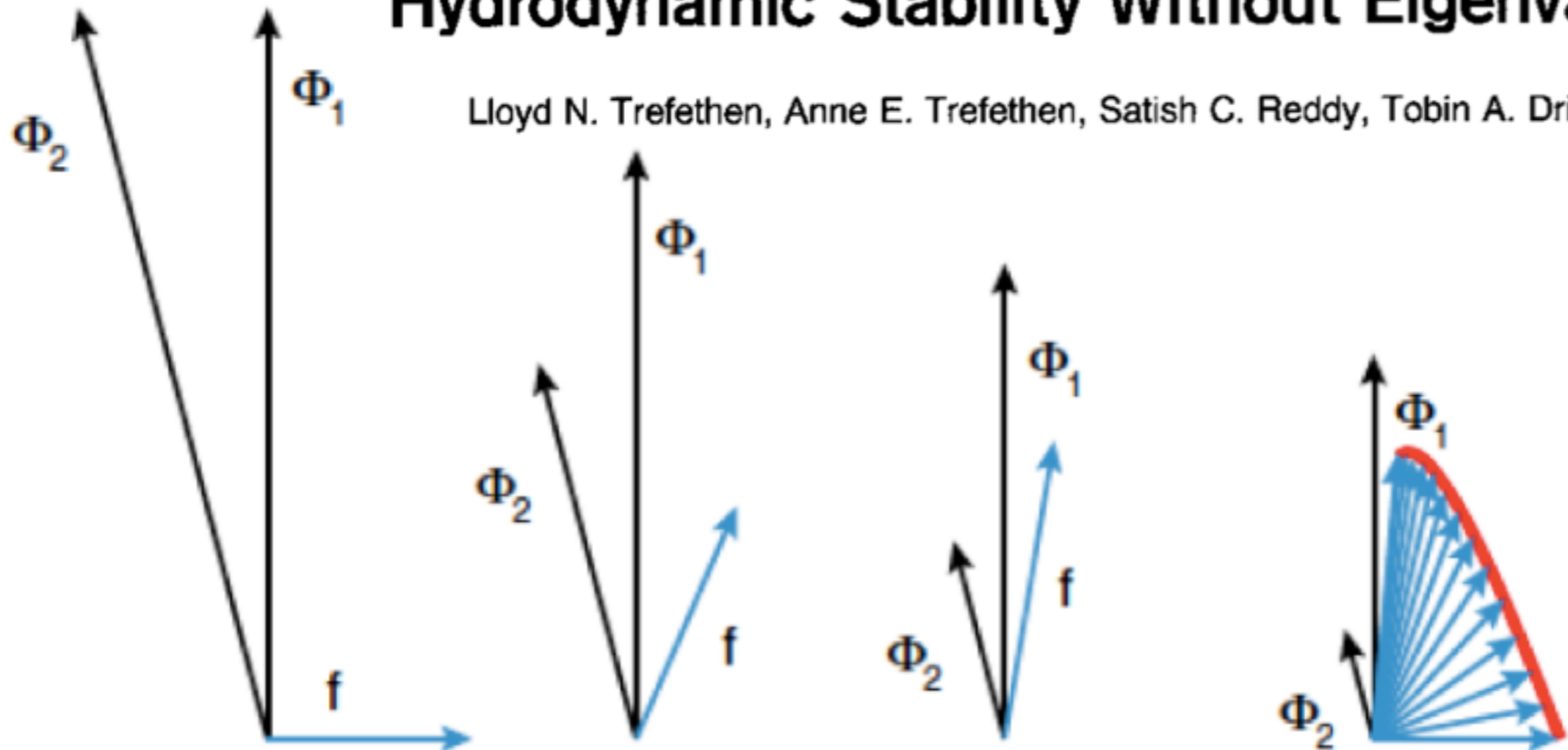
Various hydro ideas

transient growth due to non-normality of eigenvectors
transient growth by orders of magnitude

before eventual decay, nonlinearity may take over before this!

Hydrodynamic Stability Without Eigenvalues

Lloyd N. Trefethen, Anne E. Trefethen, Satish C. Reddy, Tobin A. Driscoll



subcritical transition to turbulence

Other hydro instabilities

driven by vertical rotation gradient $\frac{\partial \Omega}{\partial z}$

angular momentum transport due to vertical convection

baroclinic instability driven by $\nabla p \times \nabla \rho$

Rossby-wave instability, ...

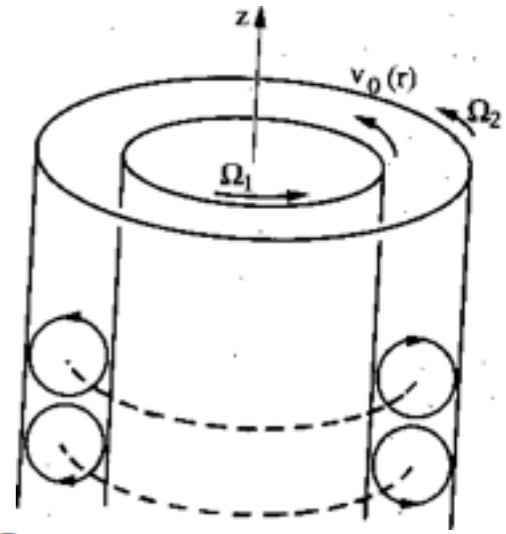
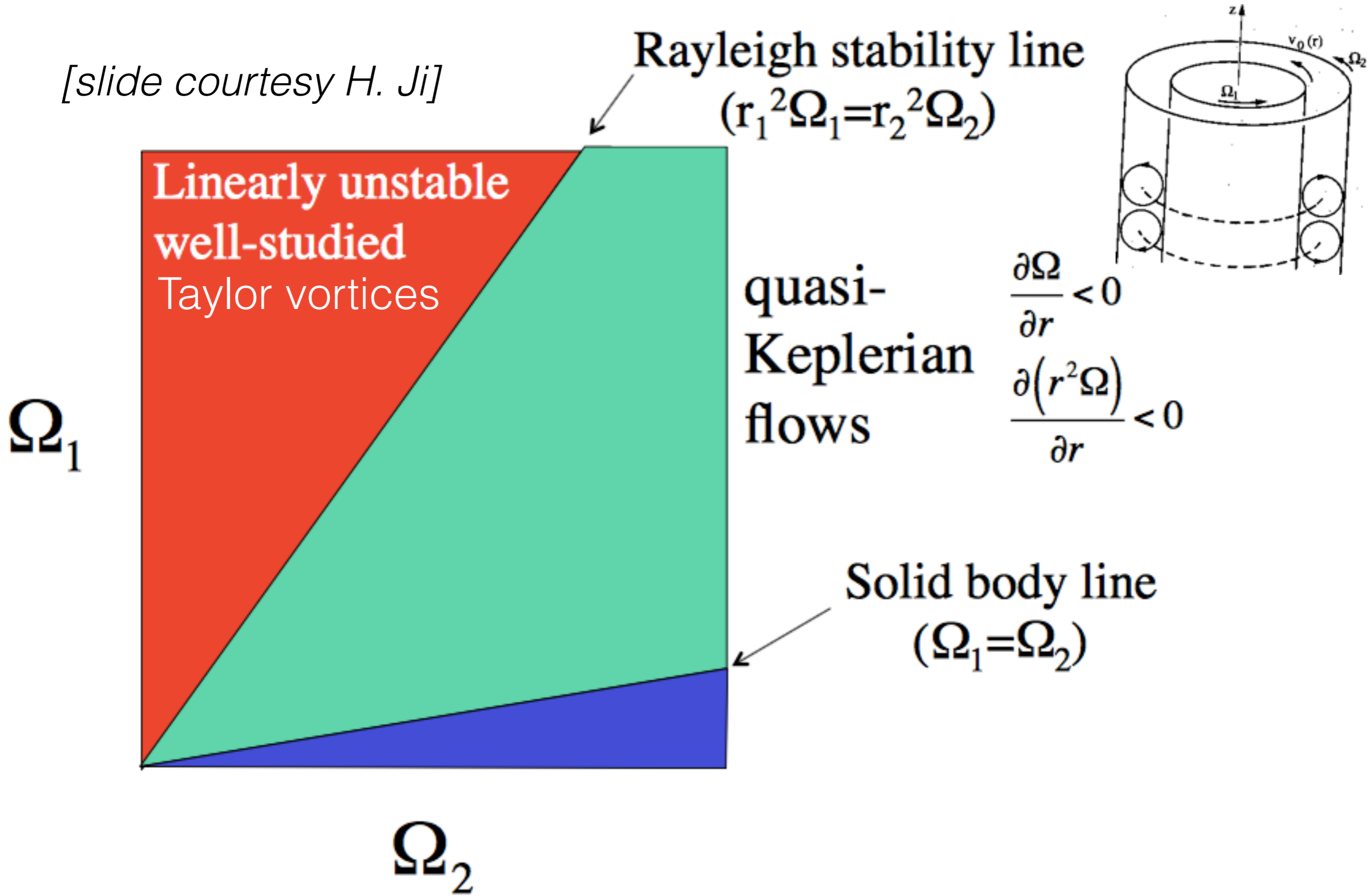
none of hydro mechanisms unanimously accepted by community

BUT PPDs must accrete: layered accretion, MHD winds?

hard problem with lot of observations!

Taylor-Couette Experiments

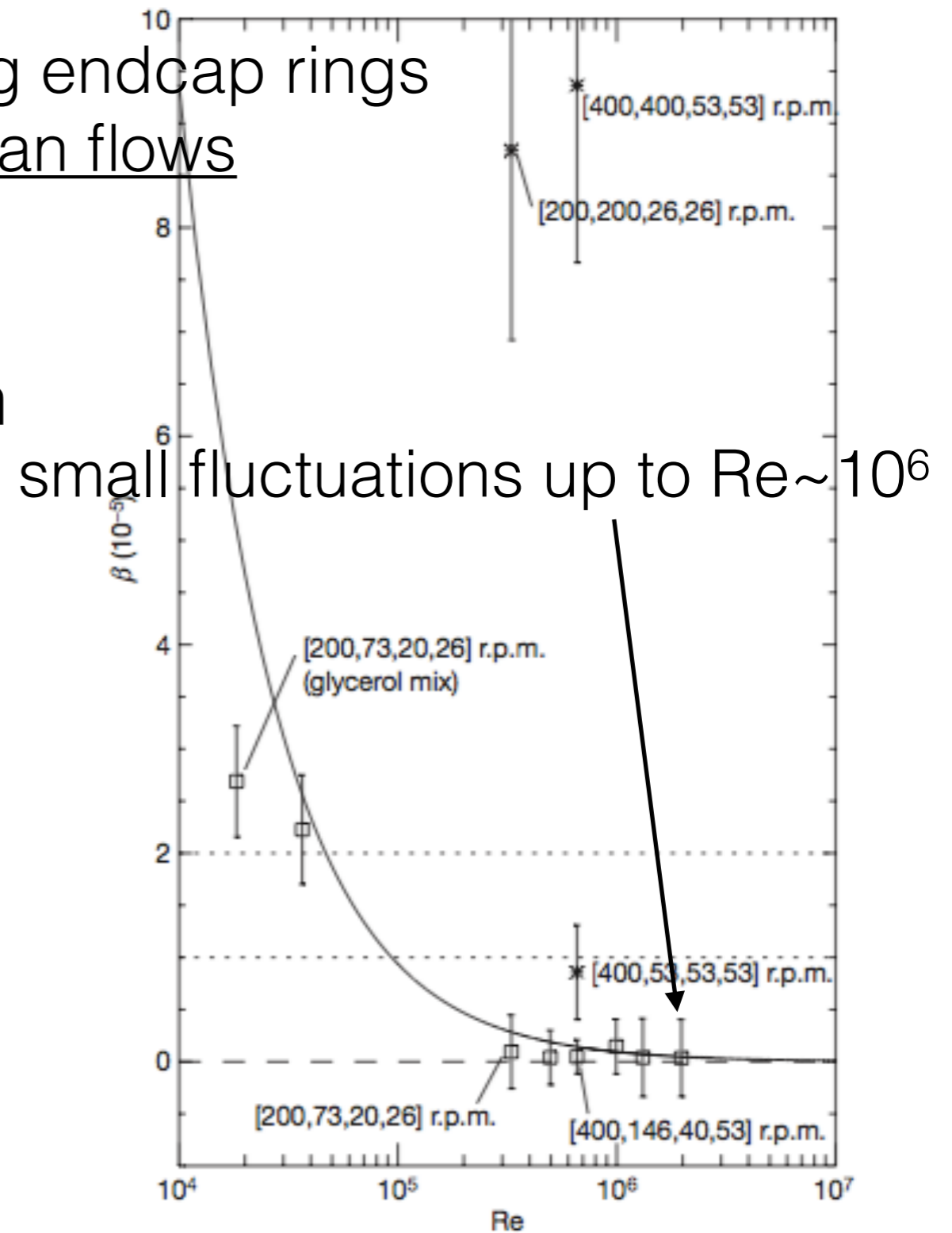
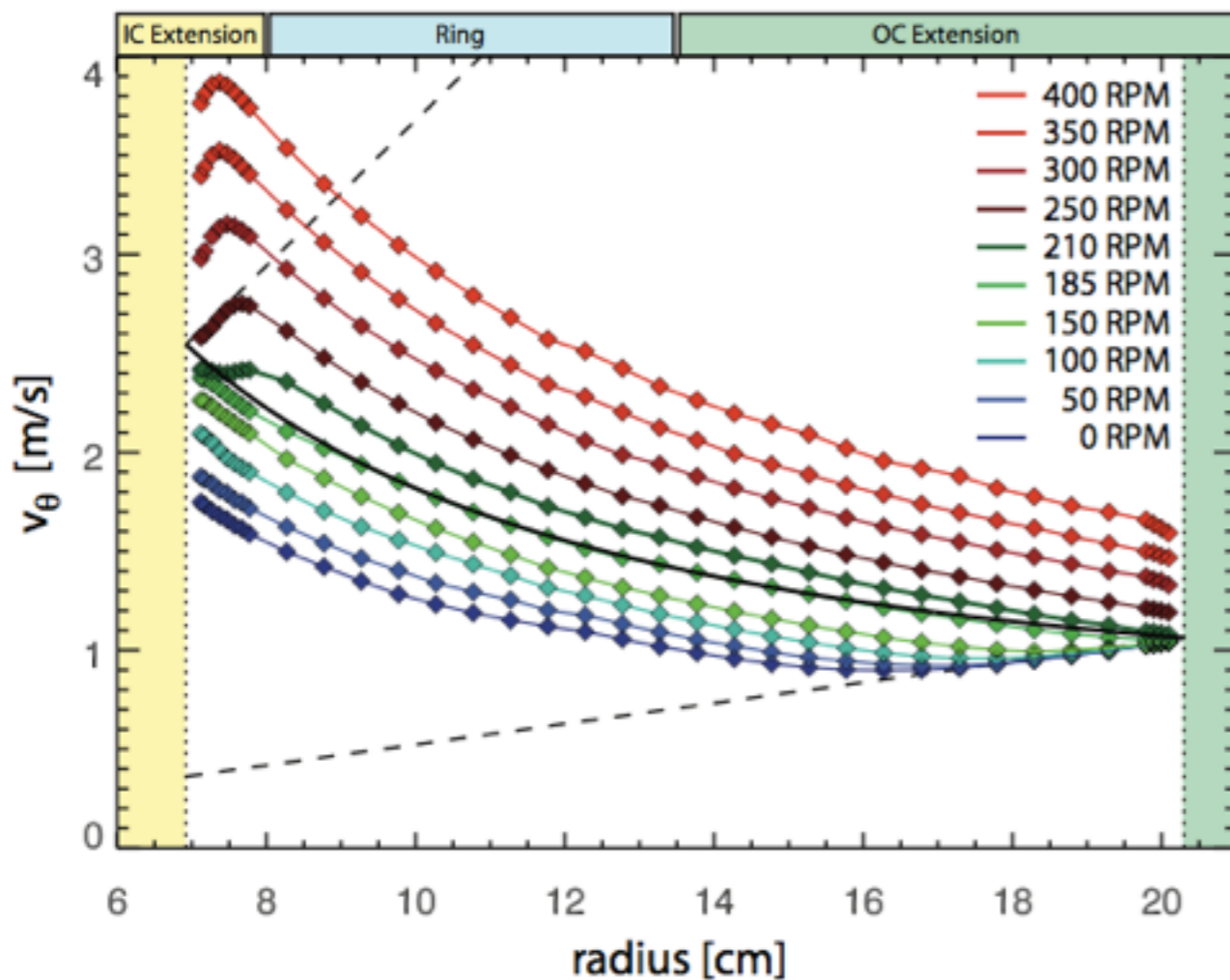
[slide courtesy H. Ji]



Experiments? [Ji et al. 2006]

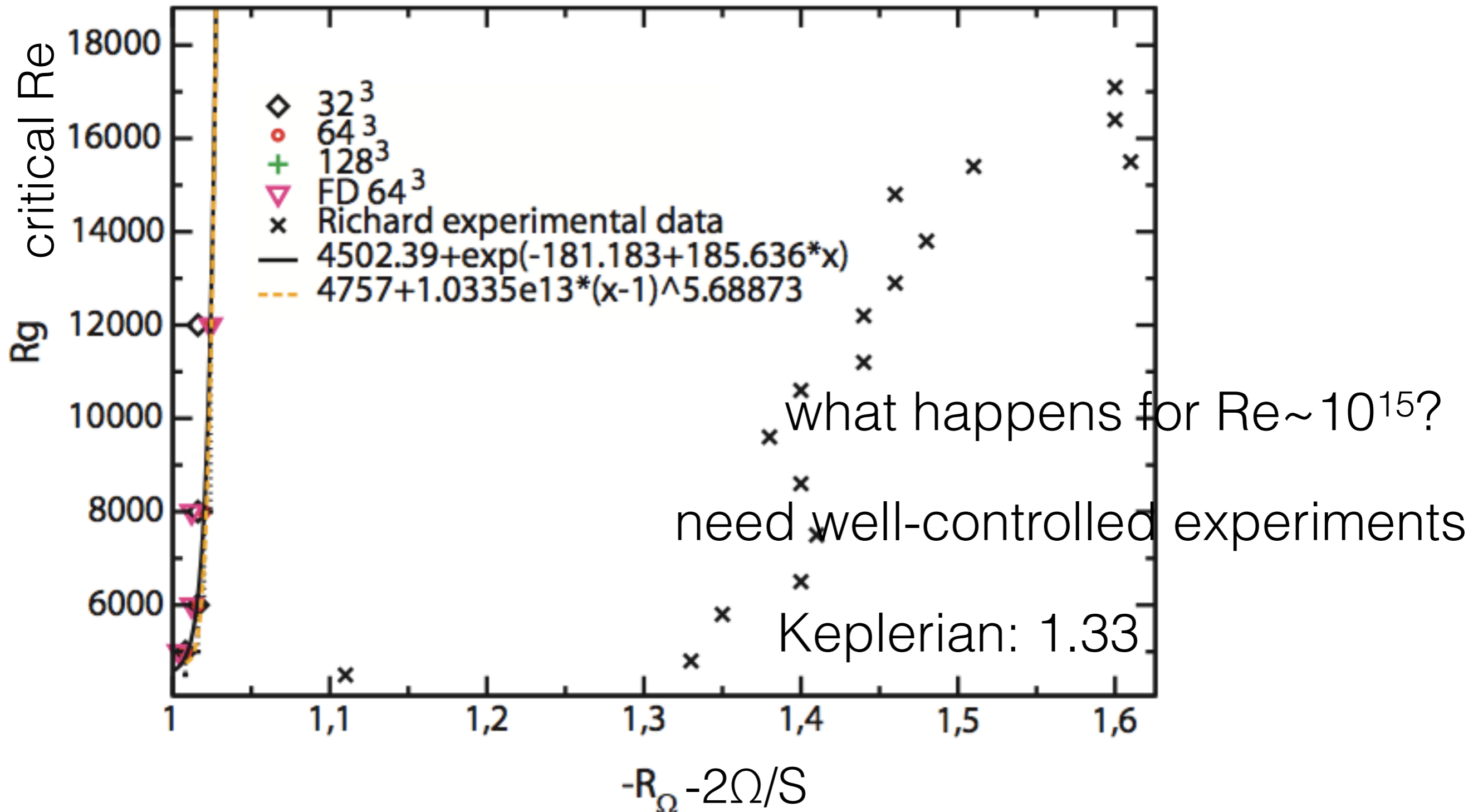
independently rotating endcap rings
to prevent Ekman flows

quasi-Keplerian rotation



Are hydro disks stable?

DNS confirm the role of Ekman boundary layers; [Avila 2012]



Summary & Future

- angular momentum transport problem
- turbulent transport: MRI mechanism for ionized flow
- what about neutral disks? linear instabilities, nonlinear mechanisms, epicyclic stabilization
- Taylor-Couette flows at $Re \sim 10^{15}$
- need experiments with controlled axial boundaries
- higher resolution local & global MHD sims with realistic microphysics (chemistry, thermodynamics, radiation, ...)
- dynamo & self sustained B-fields; role of large scale fields & outflows

Thank you